

**Relation among  
low atmospheric reconnection,  
shock formation and chromospheric jets**

**Shinsuke Takasao,  
Hiroaki Isobe and Kazunari Shibata (Kyoto Univ.)**

# Chromospheric Jets of Various Size

## Ubiquitous Reconnection

Examples:

- ▶ Spicules
- ▶ Surges
- ▶ Ca anemone jets

Ca anemone jets

time scale:

~a few-10 min.

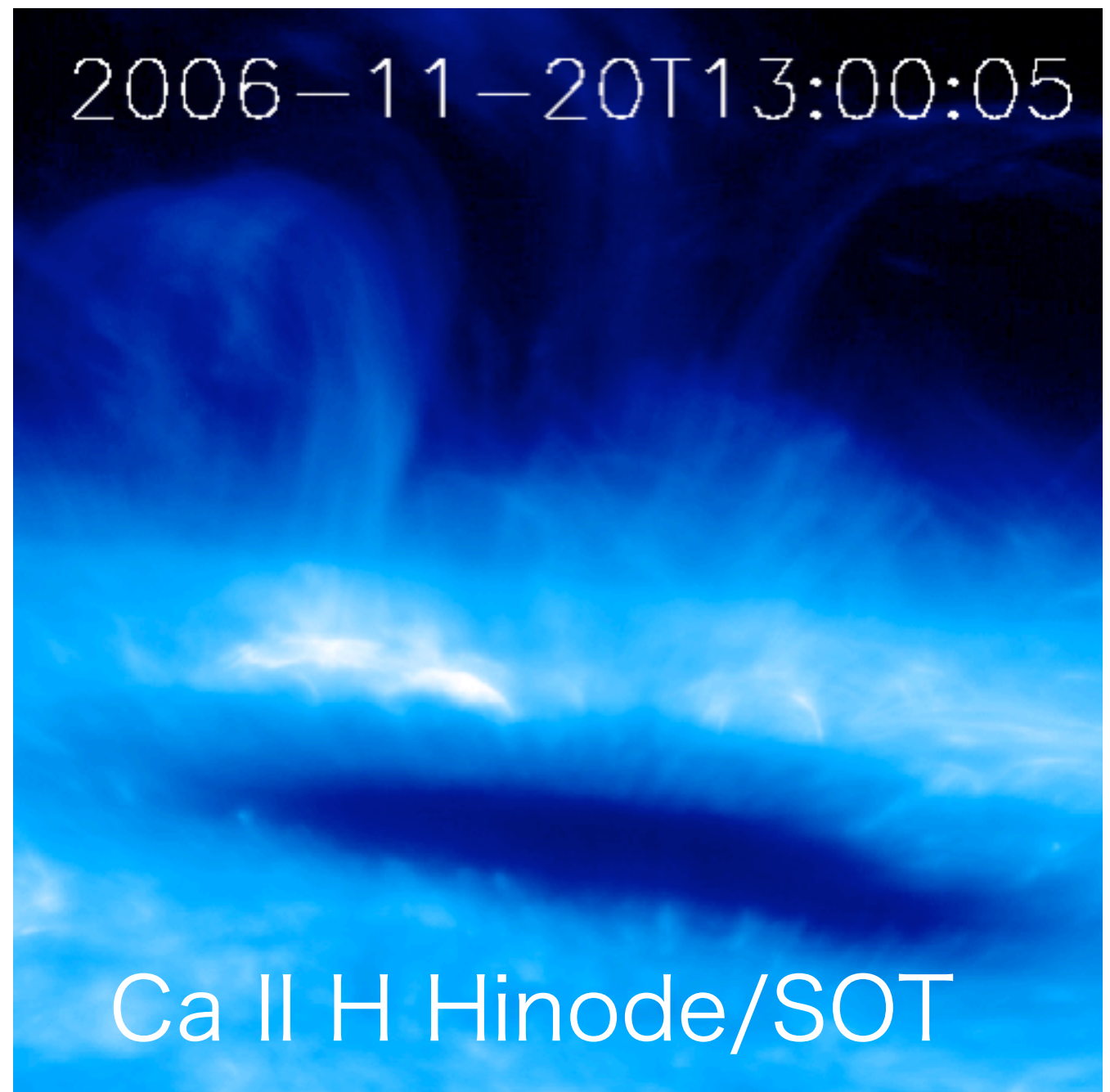
length scale:

$10^3 - 10^4$  km  $\gg$   $H_p \sim 150$  km

velocity:

5 - 20 km/s

(Nishizuka+2011)



e.g. Sterling + 1993, Shibata + 2007  
De Pontieu + 2007, Nishizuka + 2008  
Yang + 2013, Kayshap + 2013

# Classification by Size of Foot-point Structures of Jets

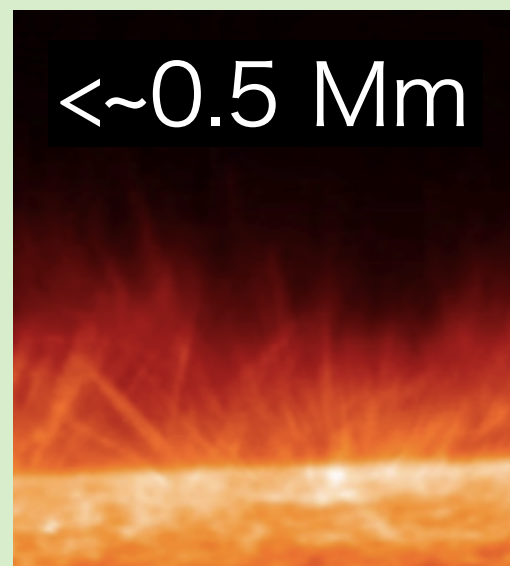
Size of foot-point structures of jets

Photo./Low chrom.

Upper chrom.

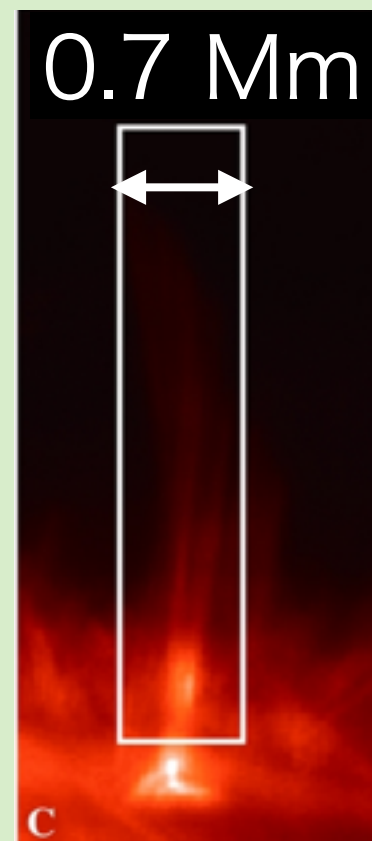
TR/corona

Spicules

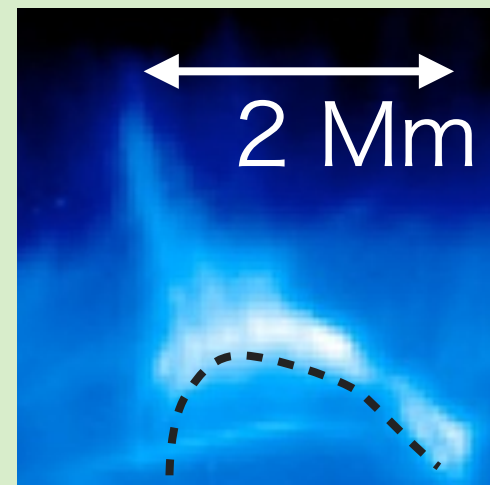


foot-points  
unresolved...

Ca jets



Singh+2012



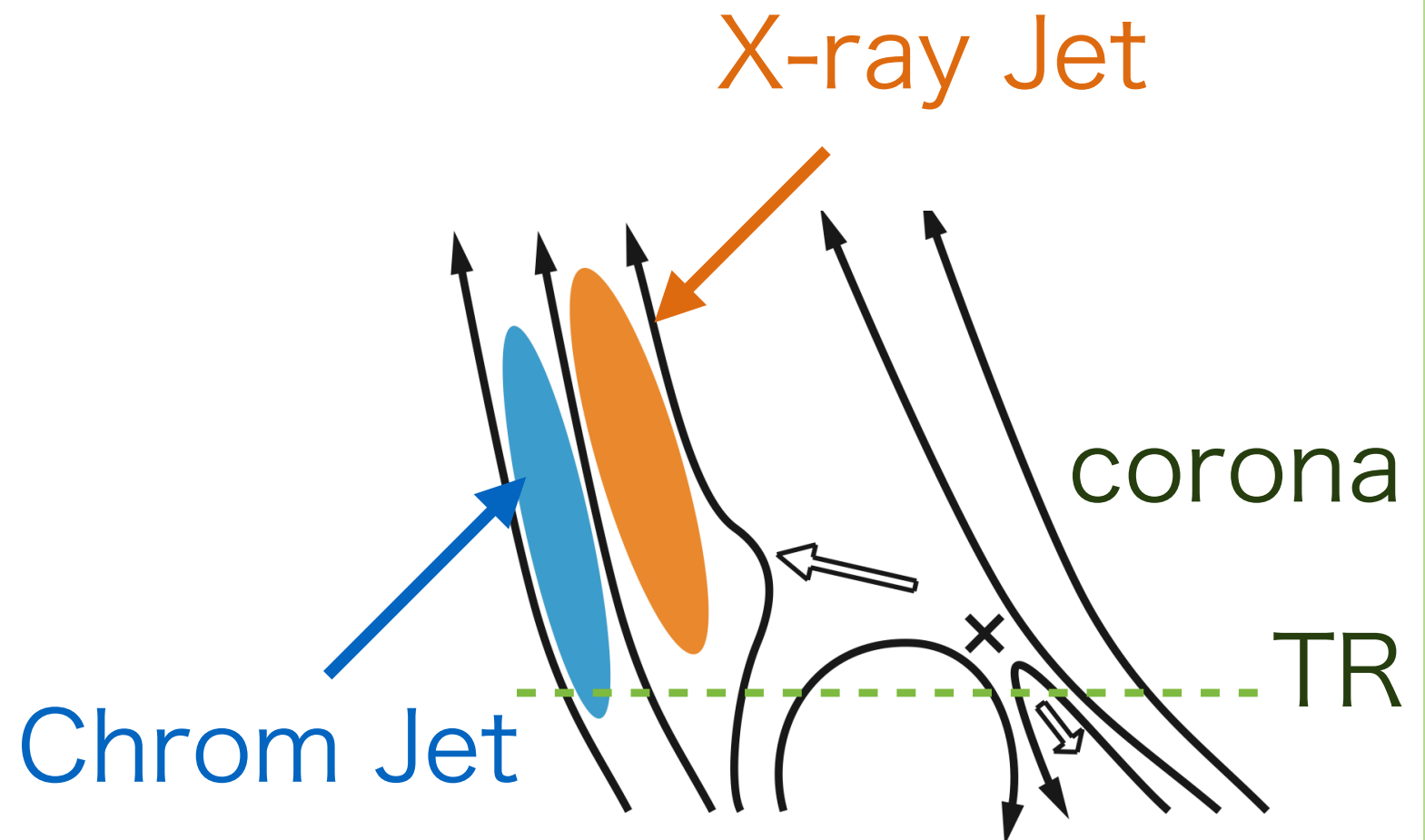
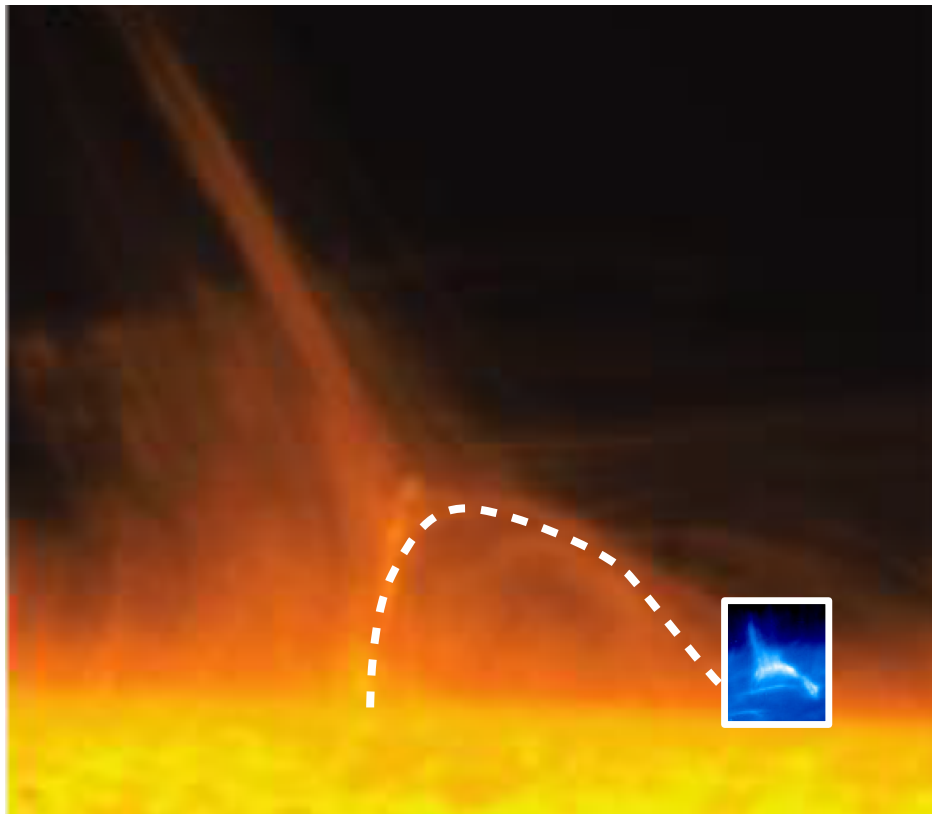
Giant Ca jets / Surges



loop- or cusp-shaped foot-points  
=> probably emerging flux

# Magnetic Reconnection at Coronal Height: Whip-like Acceleration

Surges, giant Ca jets

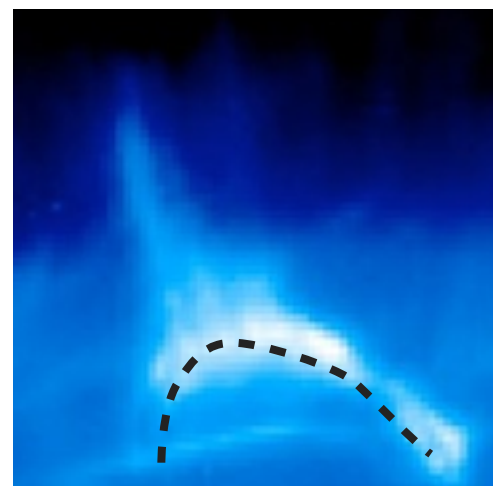
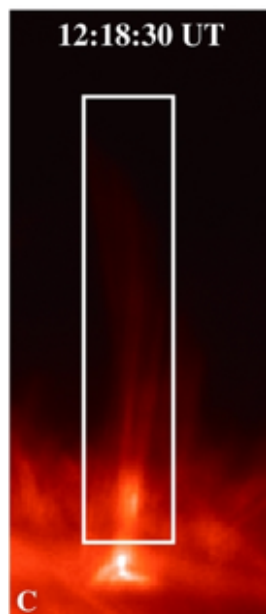
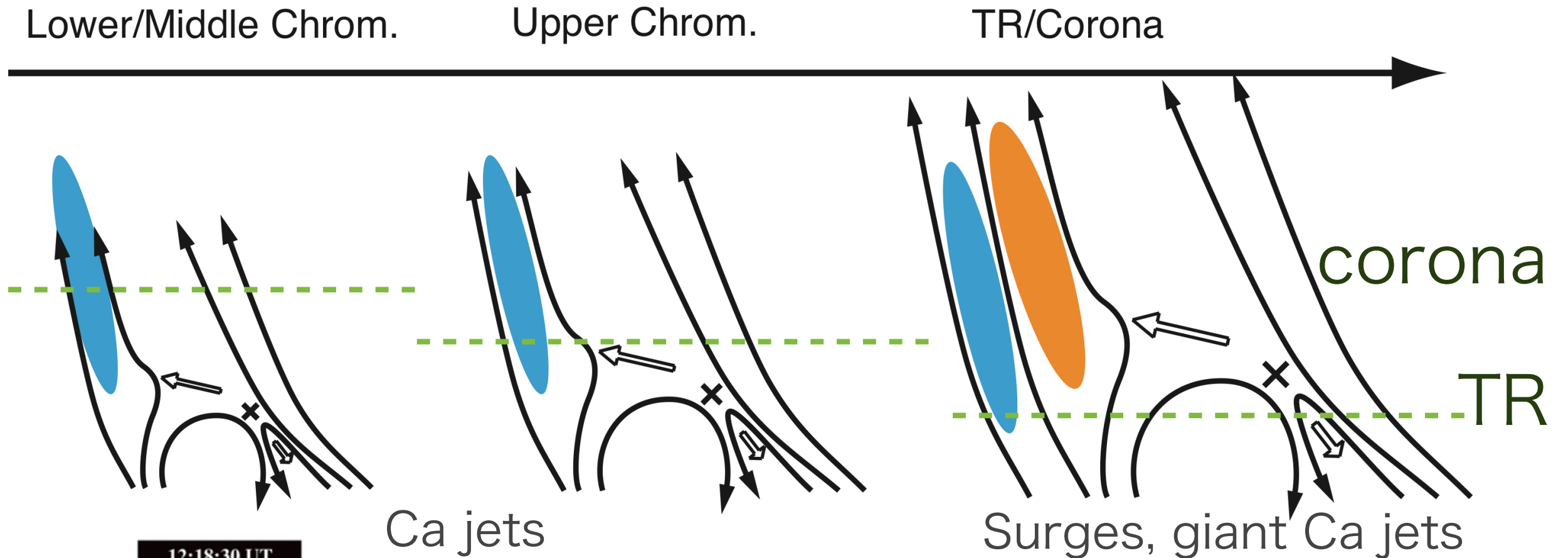


TR: Transition Region  
between chrom. and corona

Chrom. jet:  
Sling-shot  
or **whip-like acceleration**  
e.g. Yokoyama&Shibata 1996,  
Moreno-Insertis+ 2013

# Magnetic Reconnection can Take Place at Various Heights

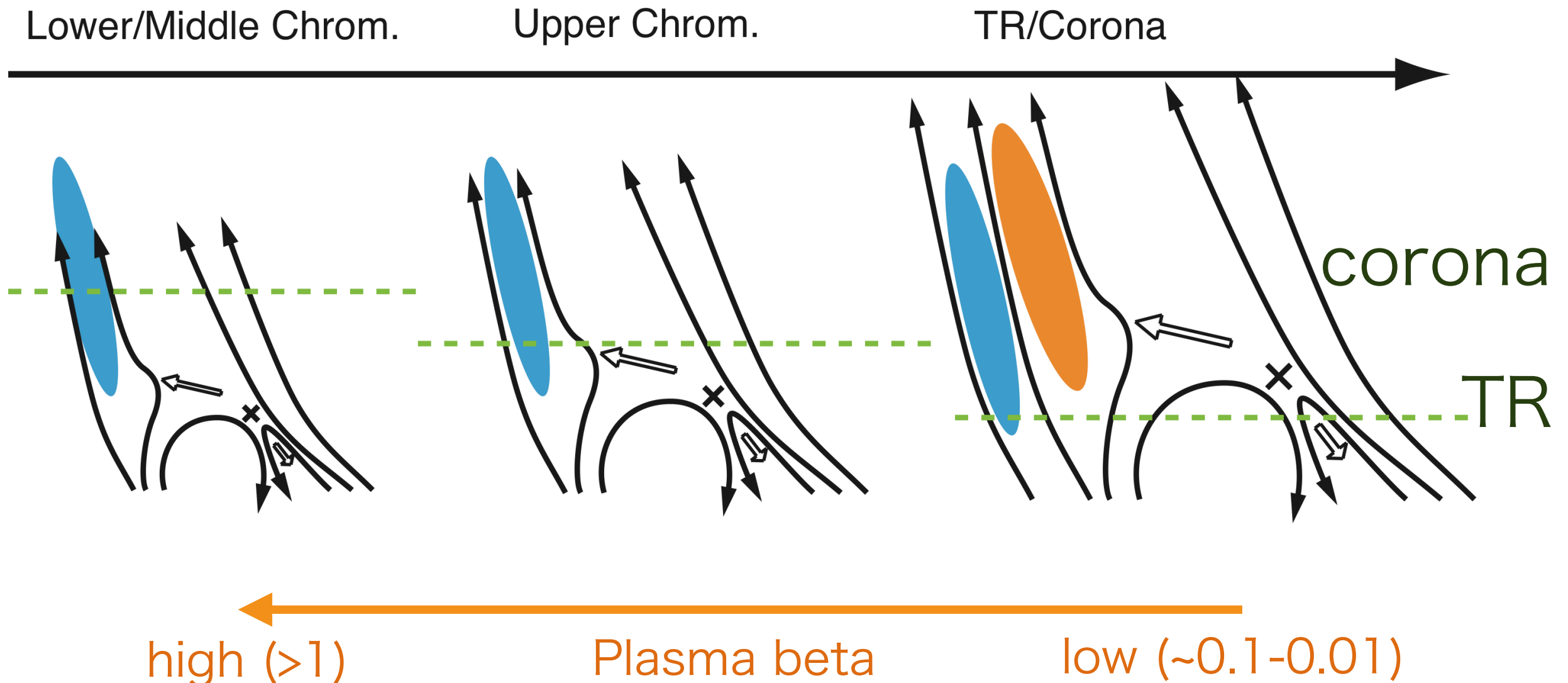
## Height of Reconnection point



Sling-shot  
or **whip-like acceleration**  
e.g. Yokoyama&Shibata 1996,  
Moreno-Insertis+ 2013

# Magnetic Reconnection can Take Place at Various Heights

## Height of Reconnection point



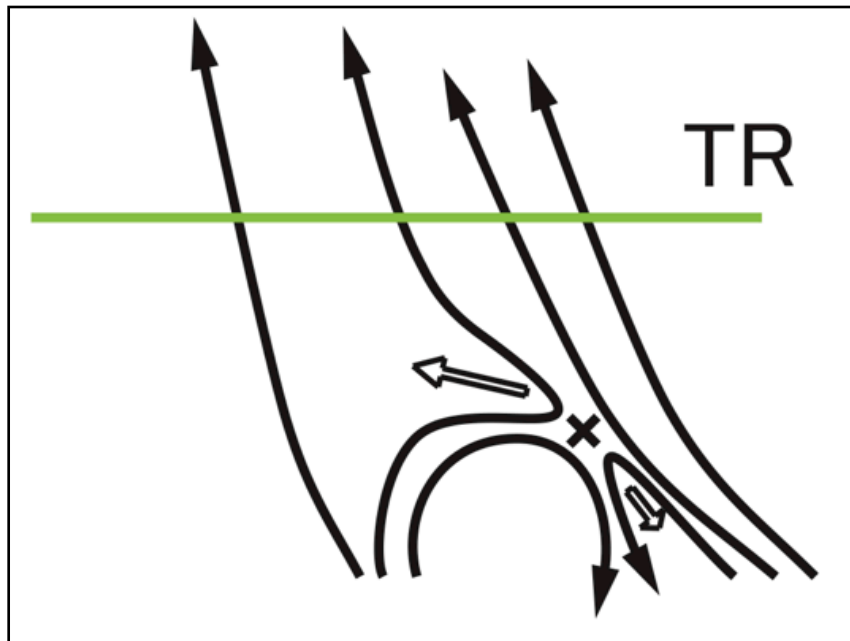
Their apparent structures are quite similar.

=> Their magnetic structures are similar.

=> Same acceleration scenario? NO!



# Low Atmospheric Reconnection and Shock



Height of jets driven by the Lorentz force  
(sling-shot / whip-like acceleration)

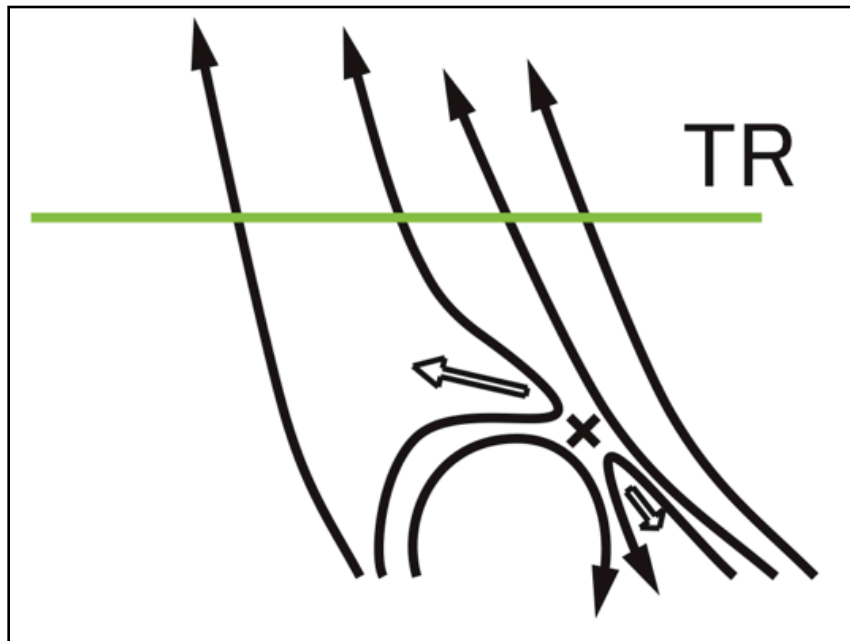
$$h_{jet} \sim H_p / \beta \sim 150 / \beta \text{ km}$$

$$h_{jet,obs} \sim 1 - 4 \times 10^3 \text{ km}$$

OK for low- $\beta$  plasma (corona)

No for high- $\beta$  plasma (low chrom.)

# Low Atmospheric Reconnection and Shock



Height of jets driven by the Lorentz force  
(sling-shot / whip-like acceleration)

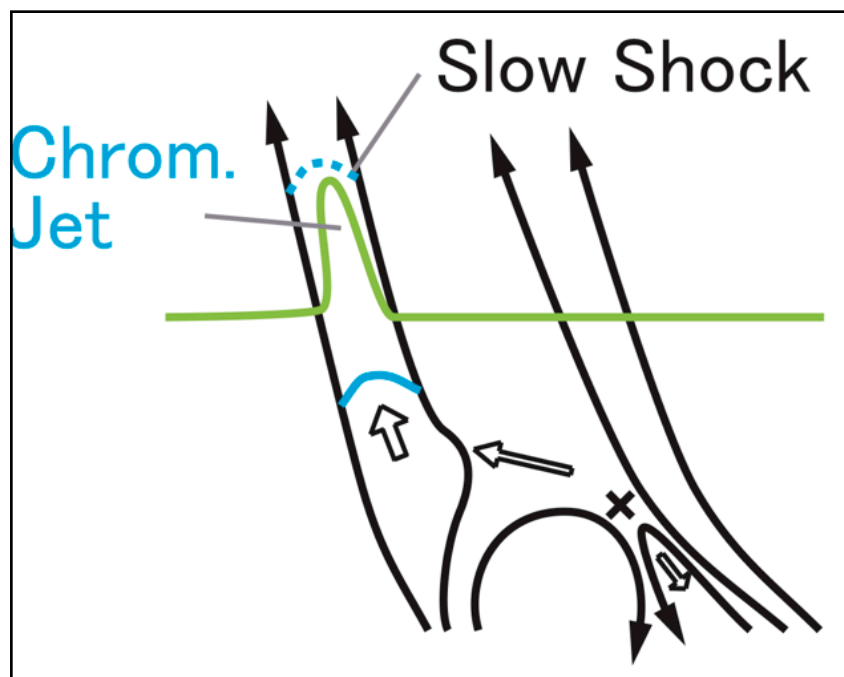
$$h_{jet} \sim H_p / \beta \sim 150 / \beta \text{ km}$$

$$h_{jet,obs} \sim 1 - 4 \times 10^3 \text{ km}$$

OK for low- $\beta$  plasma (corona)

No for high- $\beta$  plasma (low chrom.)

=> Shock acceleration



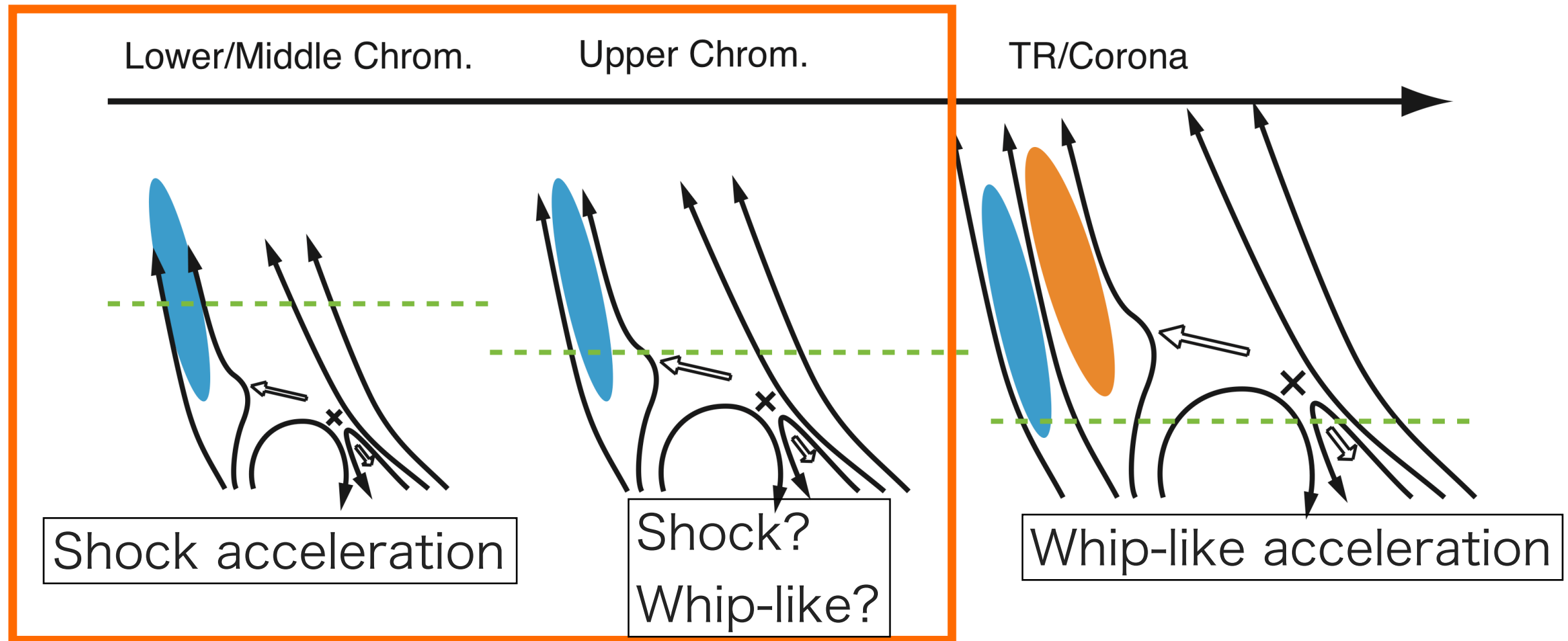
1. Energy release in the low chrom.
2. Slow mode waves/shocks carry the energy along a magnetic field
3. Only a fraction of the plasma in the upper chrom. (low-density plasma) is accelerated by shocks.

(e.g. Shibata+1982, 2007,  
Hegglund+2007)

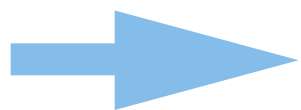


# Acceleration Mechanism Depends on Height of Reconnection

Height of Reconnection point

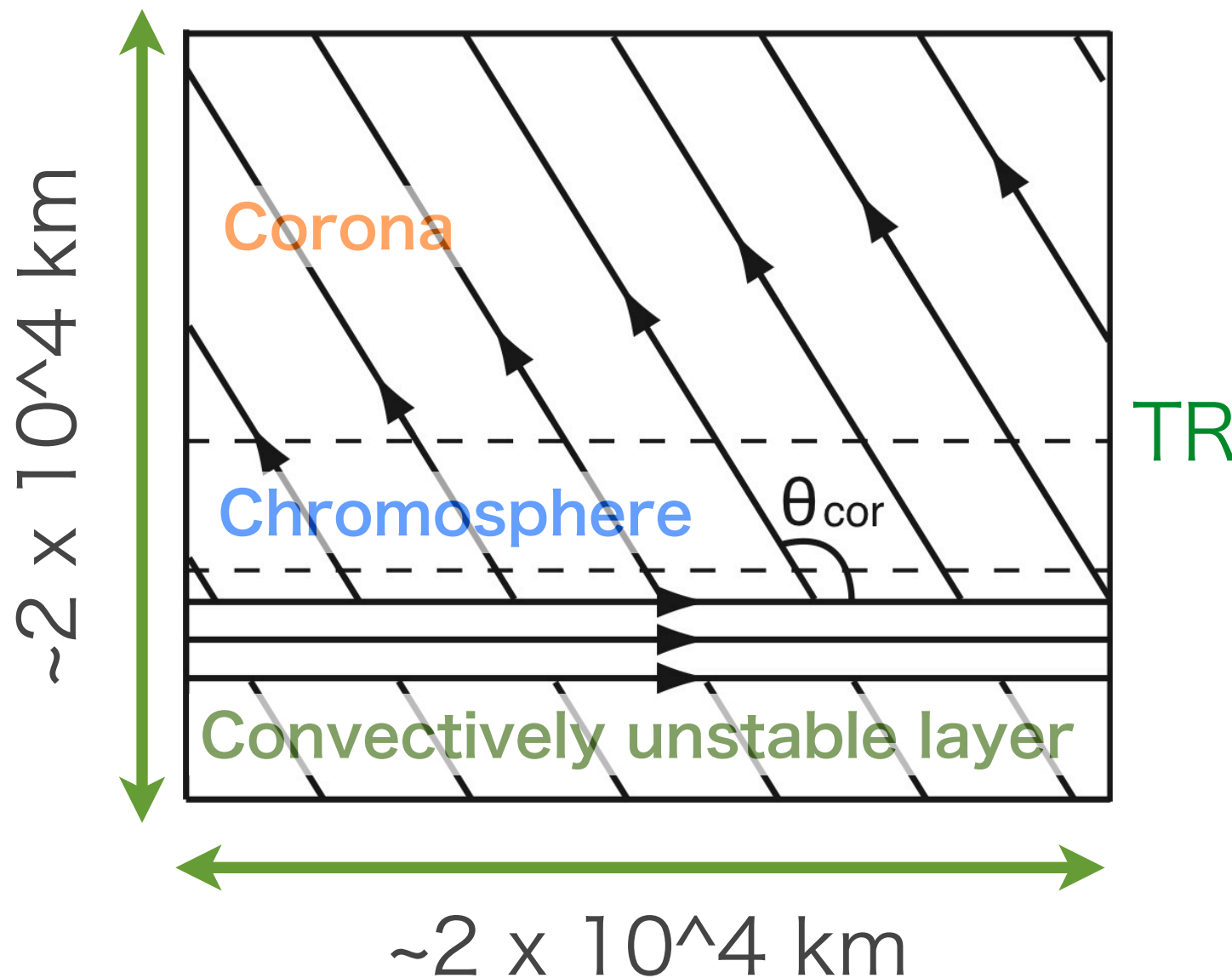


How shocks form? How shocks accelerate chrom. jets?



MHD simulation

# Initial Conditions and Assumptions of Simulation

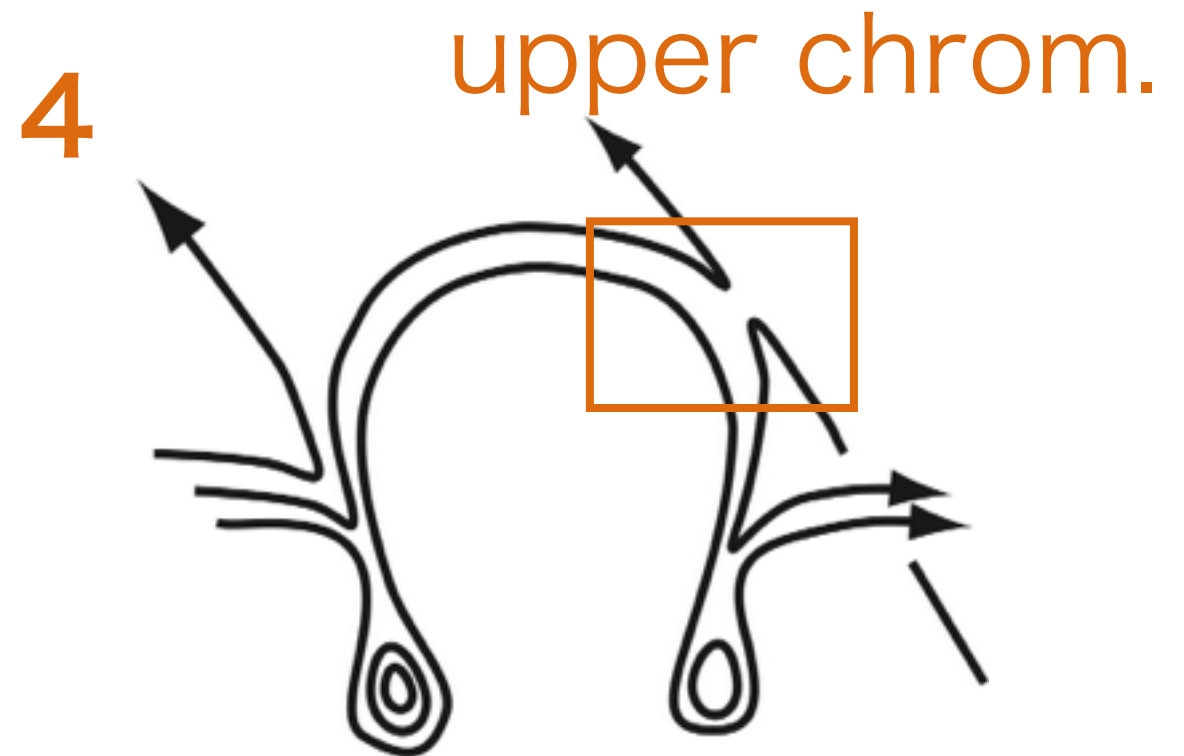
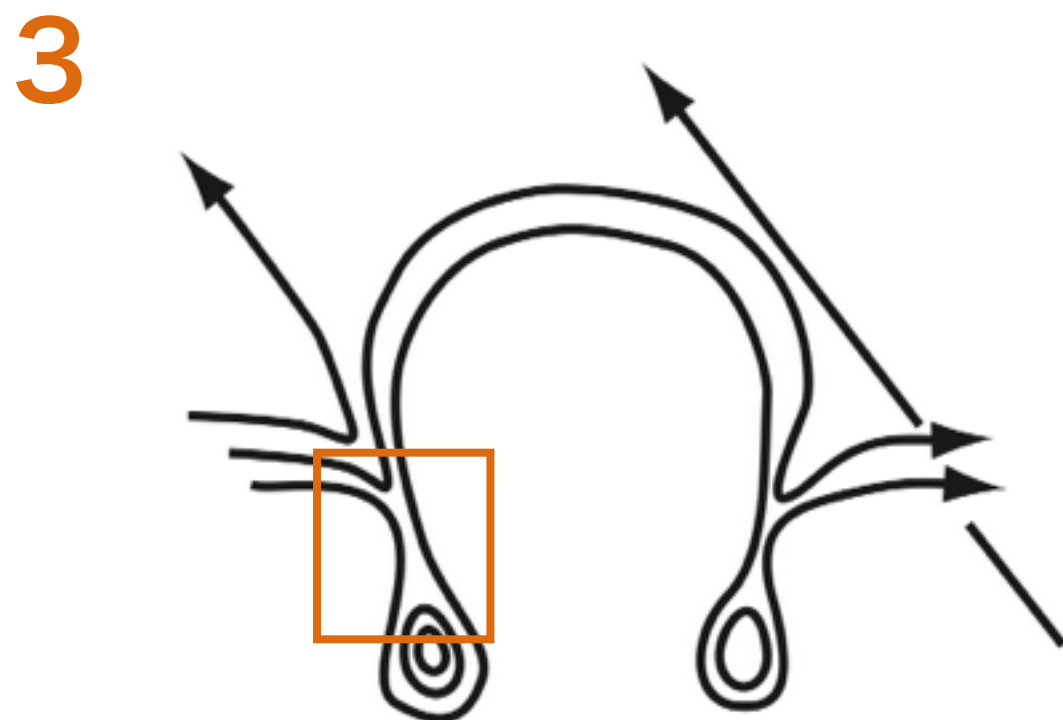
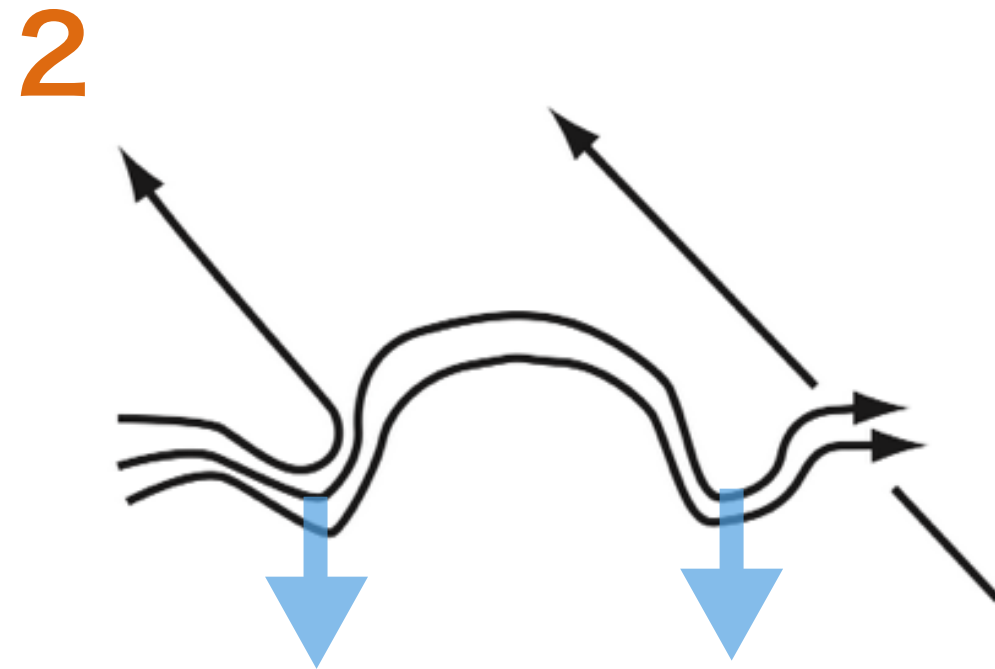
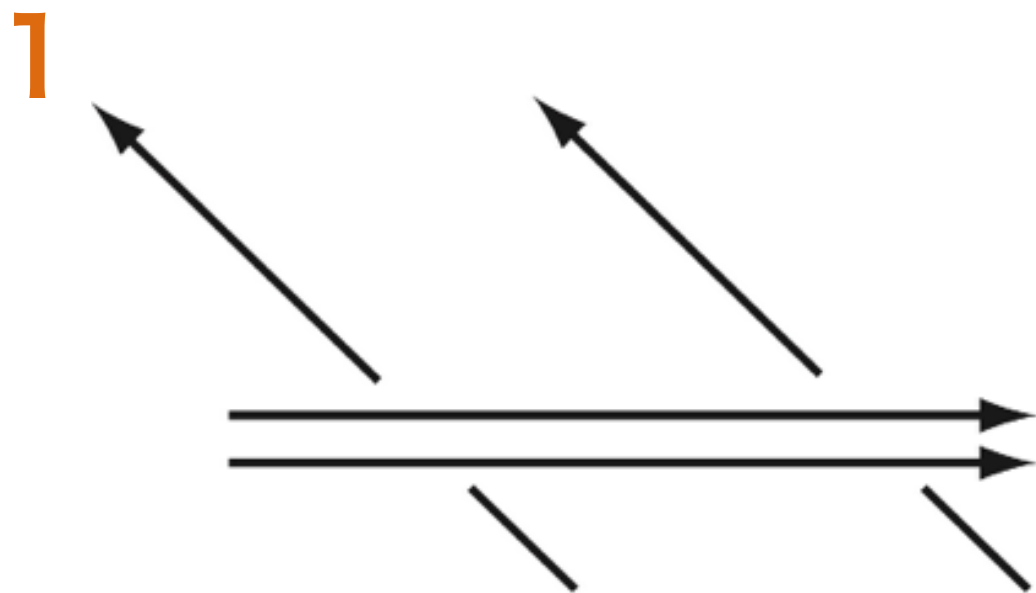


Assumptions:

- uniform gravity
- 2D MHD
- simple radiation cooling func. in the chrom.
- $\theta_{\text{cor}} = 2\pi/3$
- Anomalous resistivity model (a localized resistivity model)

TR: transition region, where the temperature / density drastically vary (density discontinuous layer)

# Where Reconnection Takes Place?

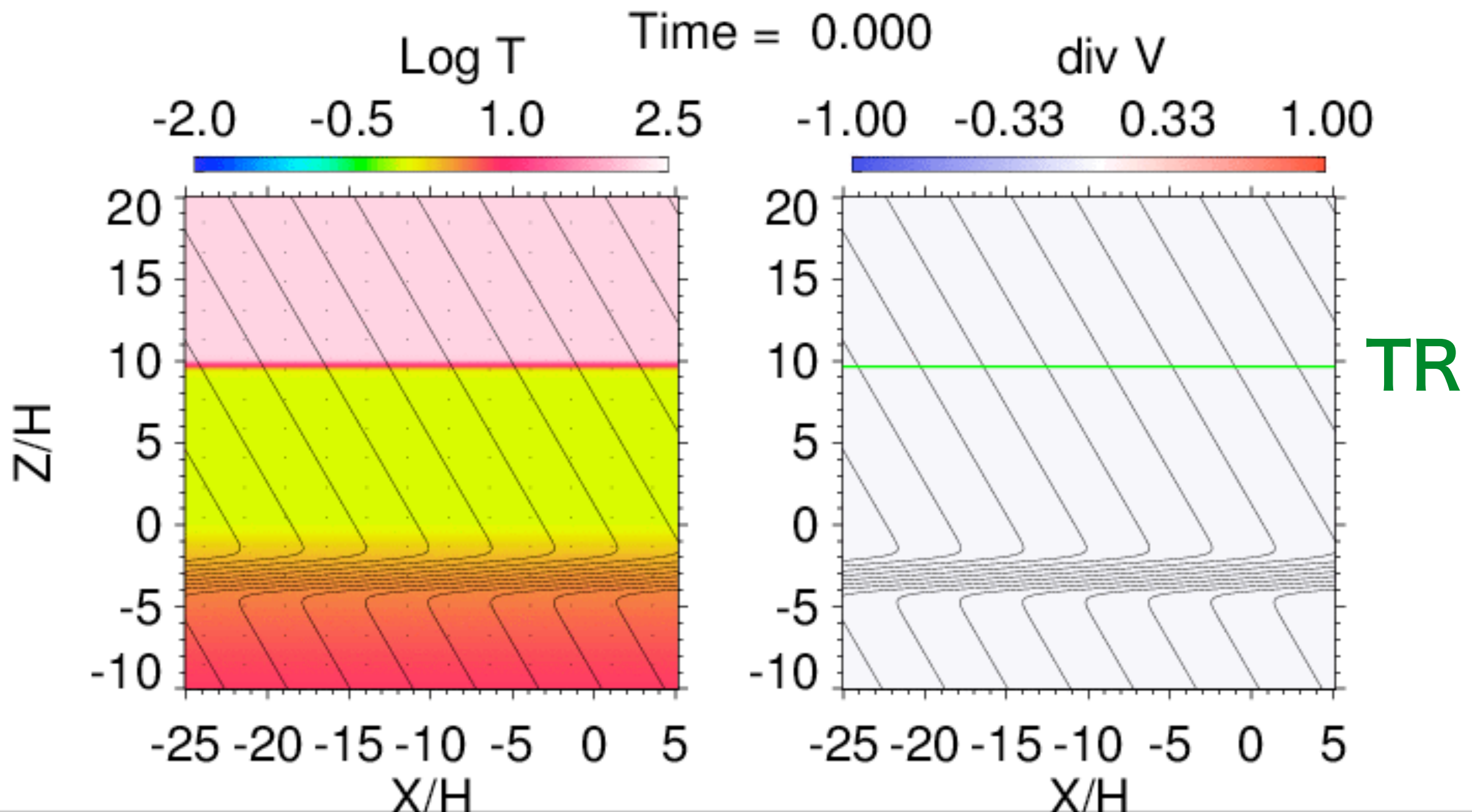


just below the photo.

# Numerical Results

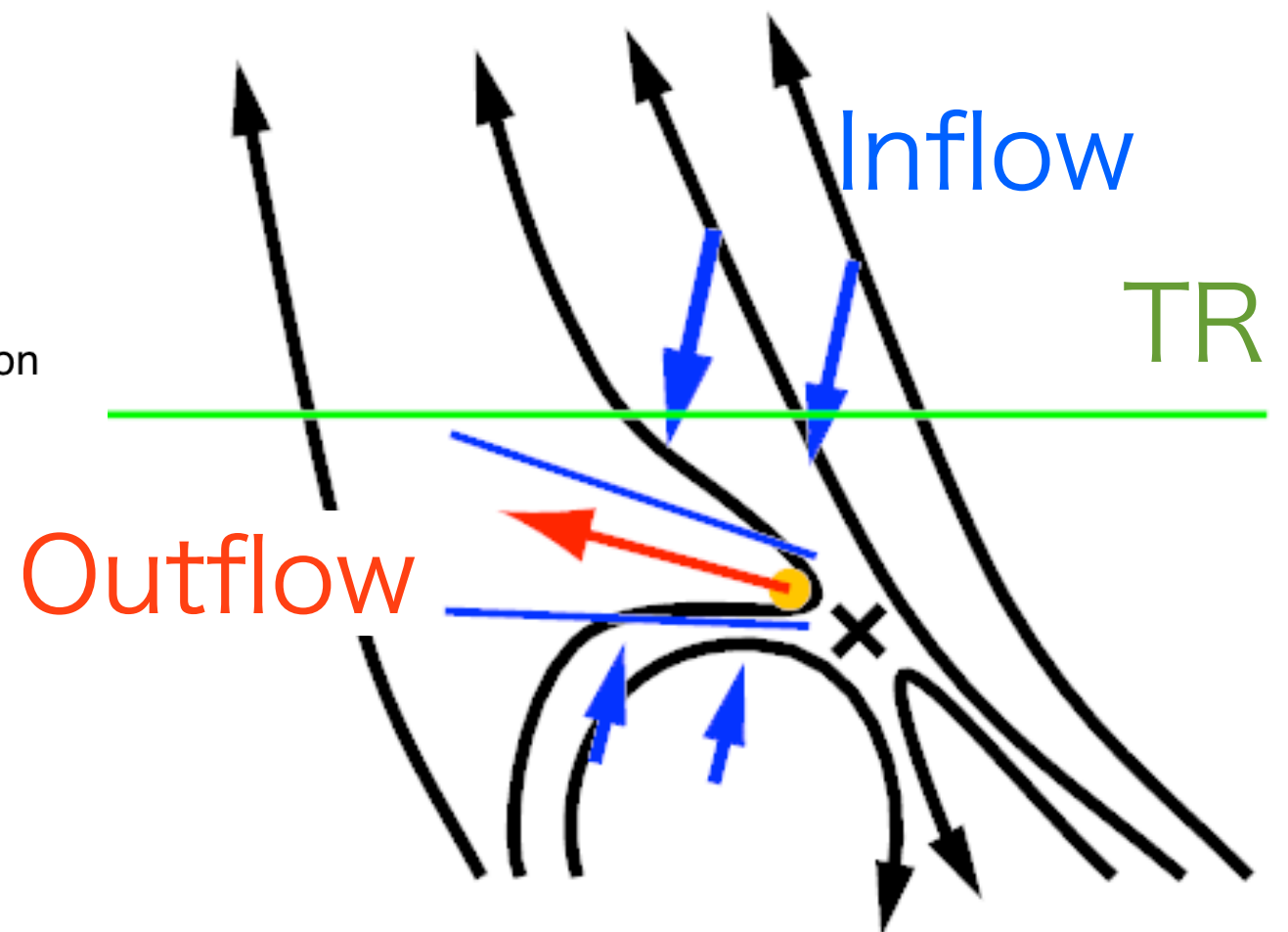
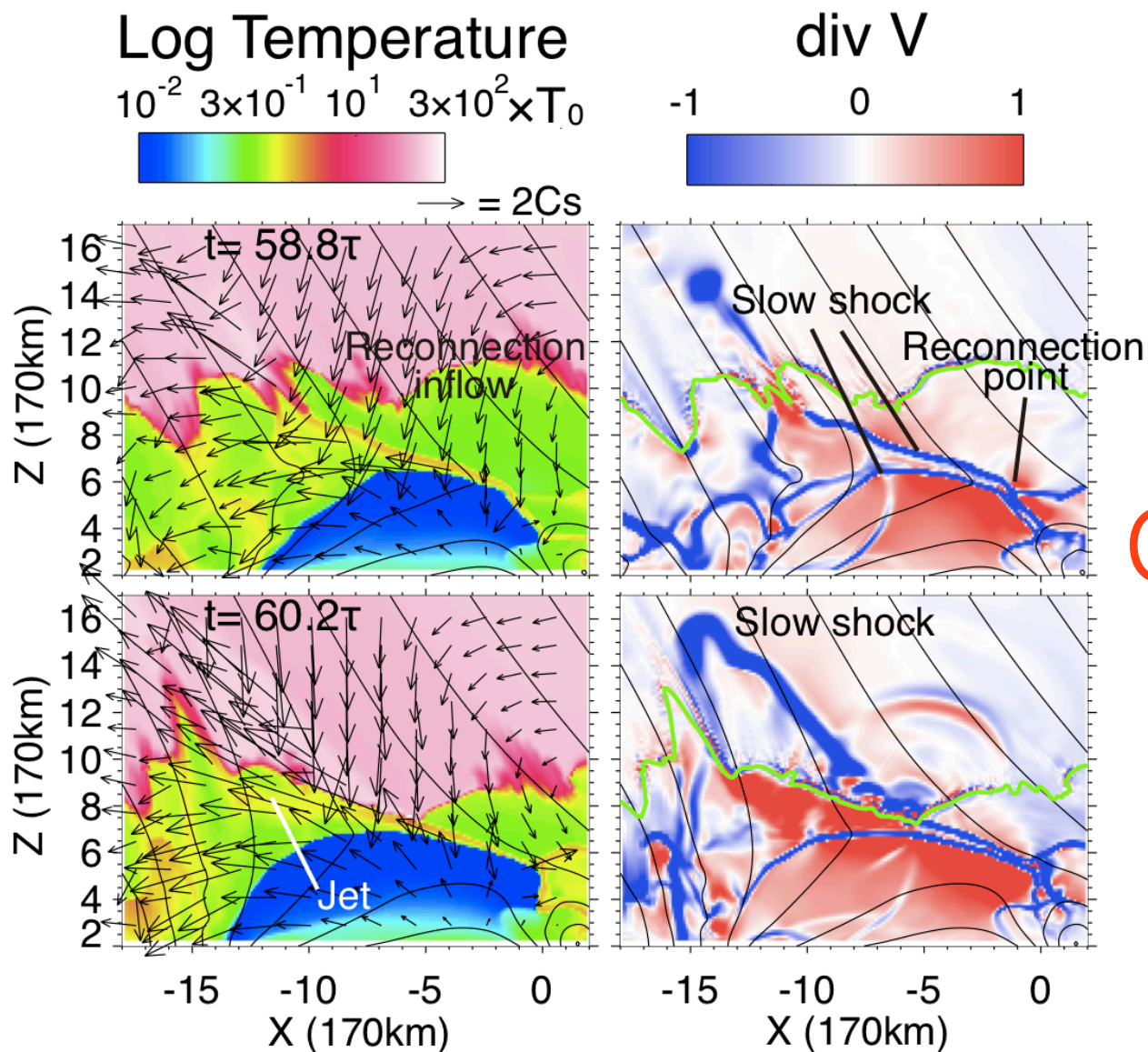
divergence  
of the velocity field

Blue : compressed (~shocks)  
Red : expanded



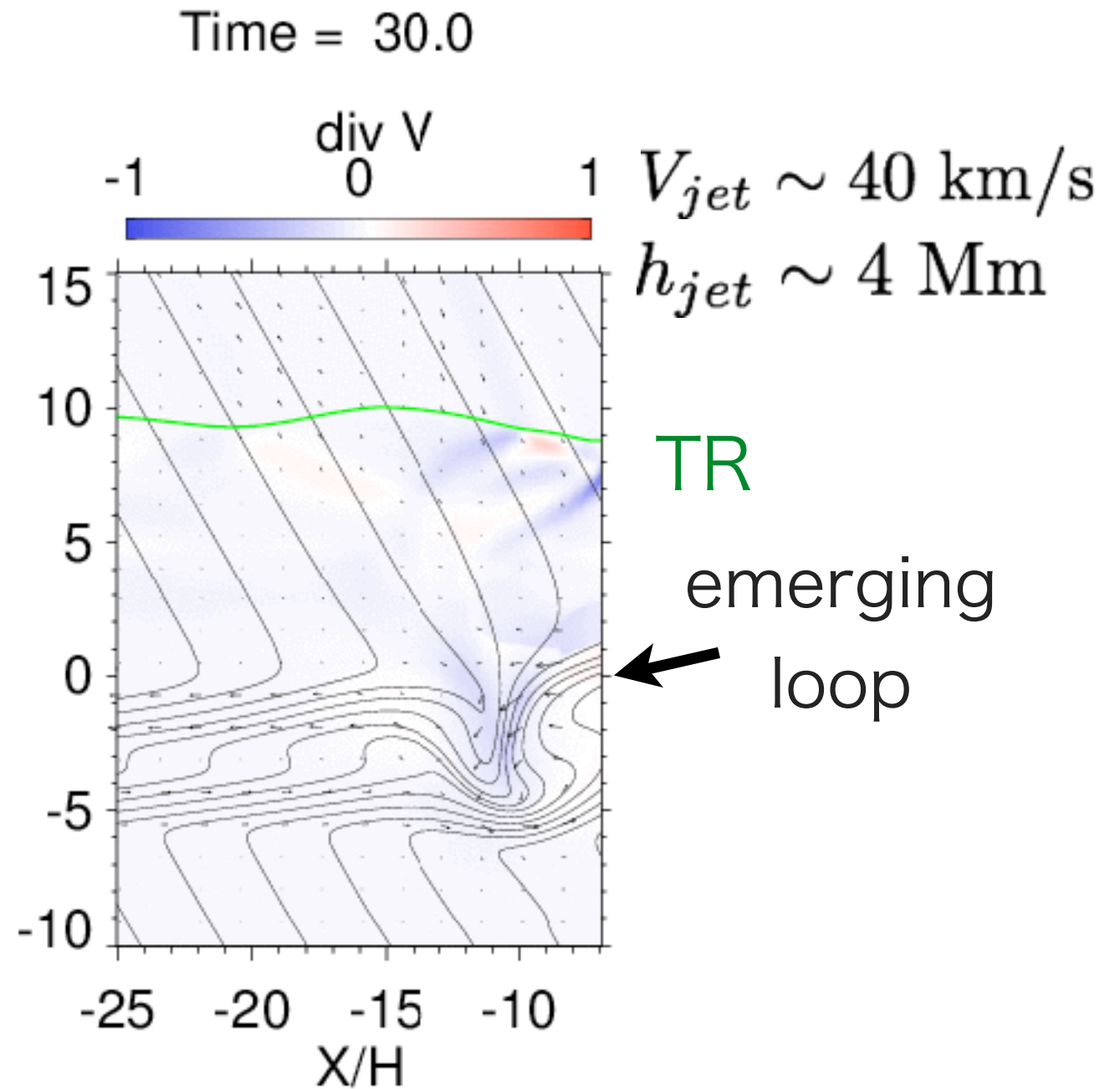
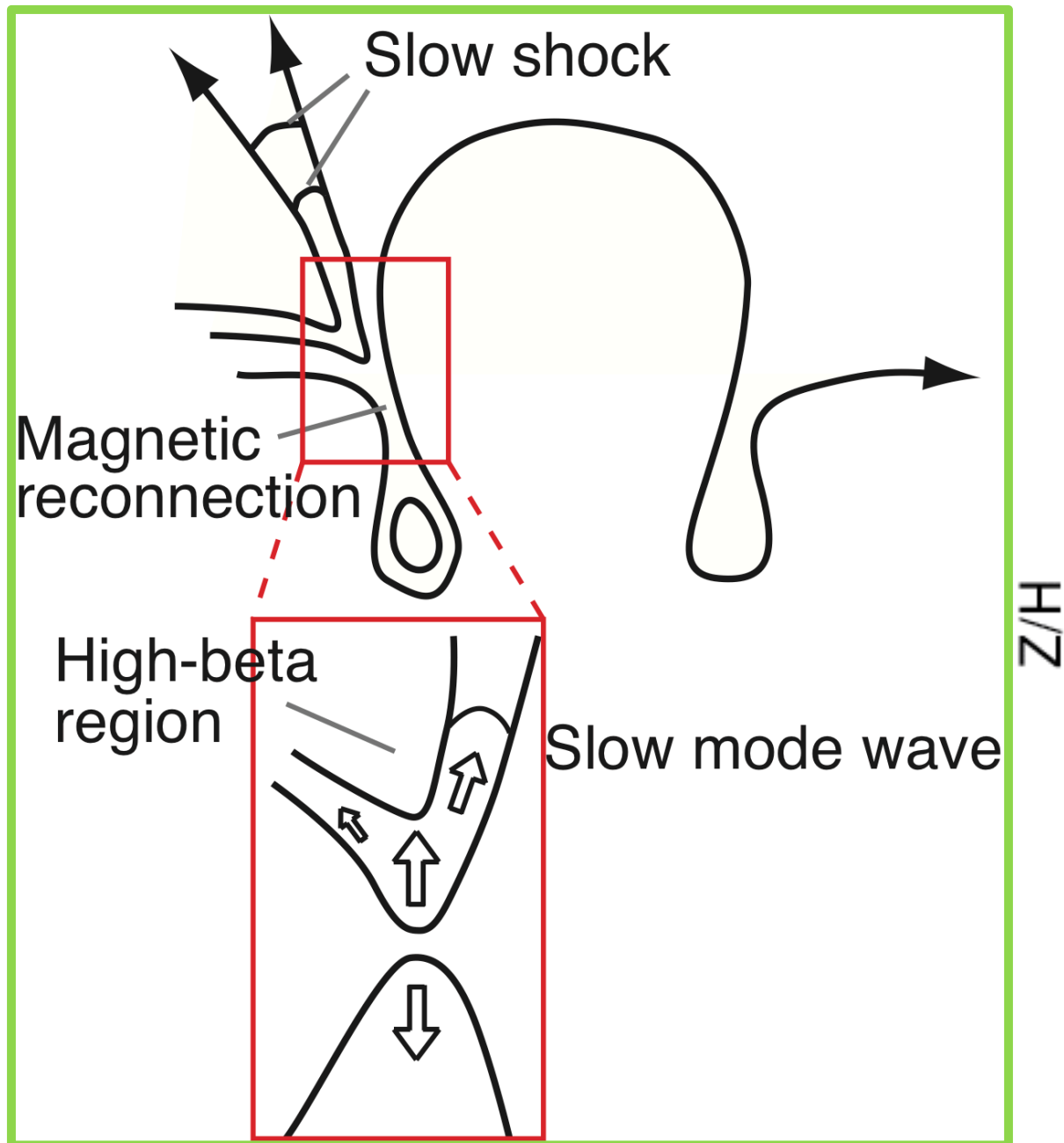


# Upper Chrom. Recon.: Combination of Whip-like and Shock Acceleration



- The chrom plasma is accelerated by
1. the magnetic tension  
(whip-like acceleration)
  2. the slow shock  
(slow shock acceleration)

# Lower Atmospheric Recon.: Shock Acceleration



Application:

Ellerman bombs => H-alpha Surges (e.g. Pariat+2004)



# Systematic Understanding of Chromospheric Jets: Classification by the Height of Recon. Points

Height of Reconnection Point

Lower/Middle Chrom.

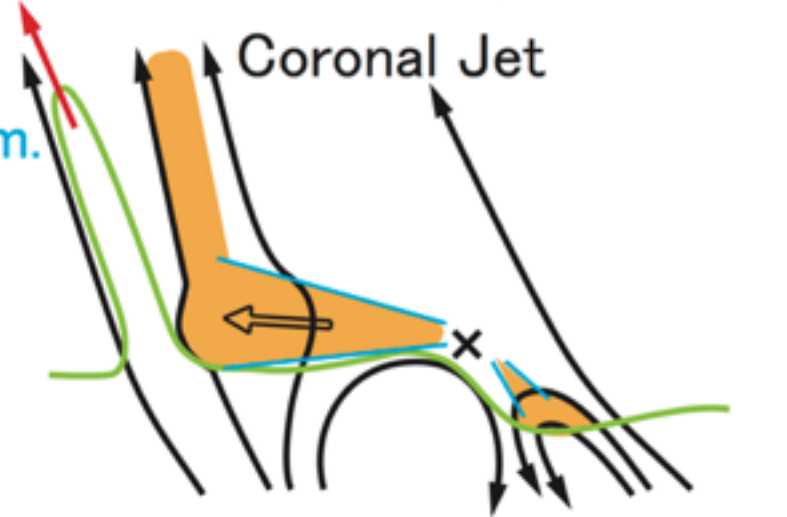
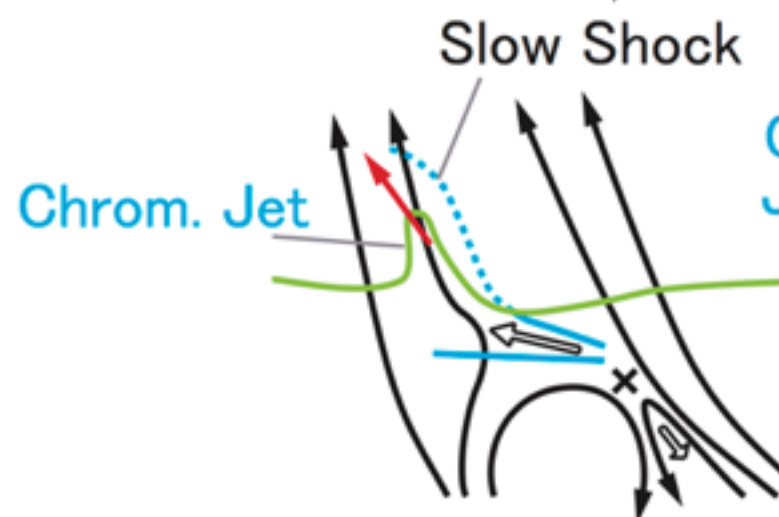
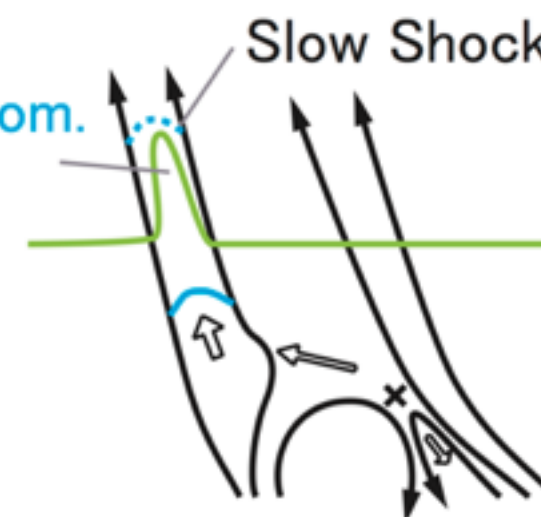
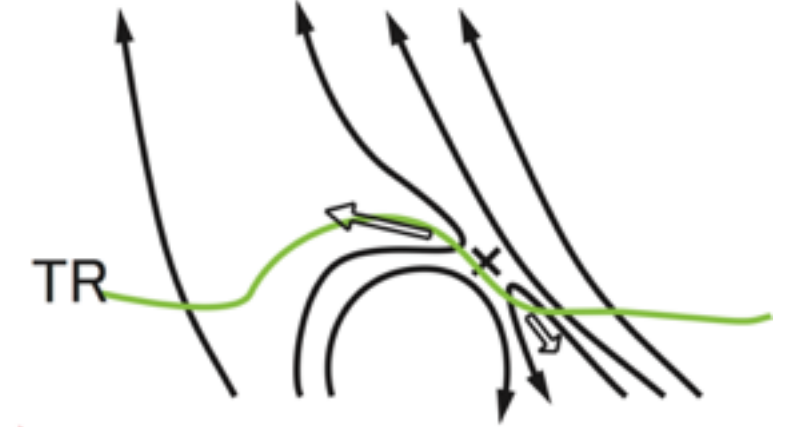
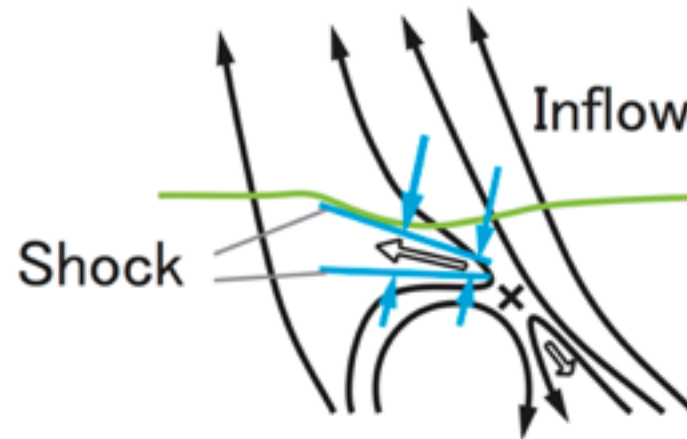
Upper Chrom.

TR/Corona →

Shock Acceleration

Shock + Whip-like Acceleration

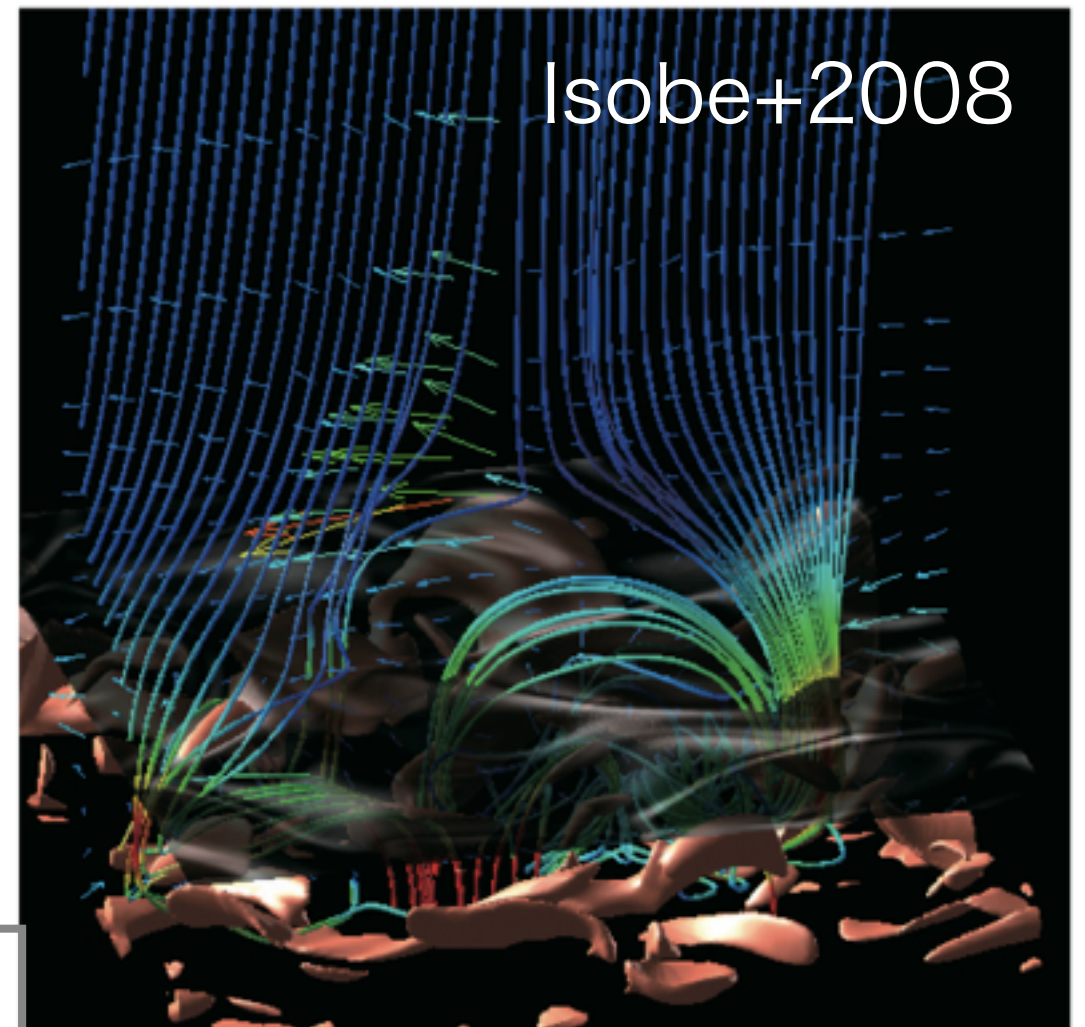
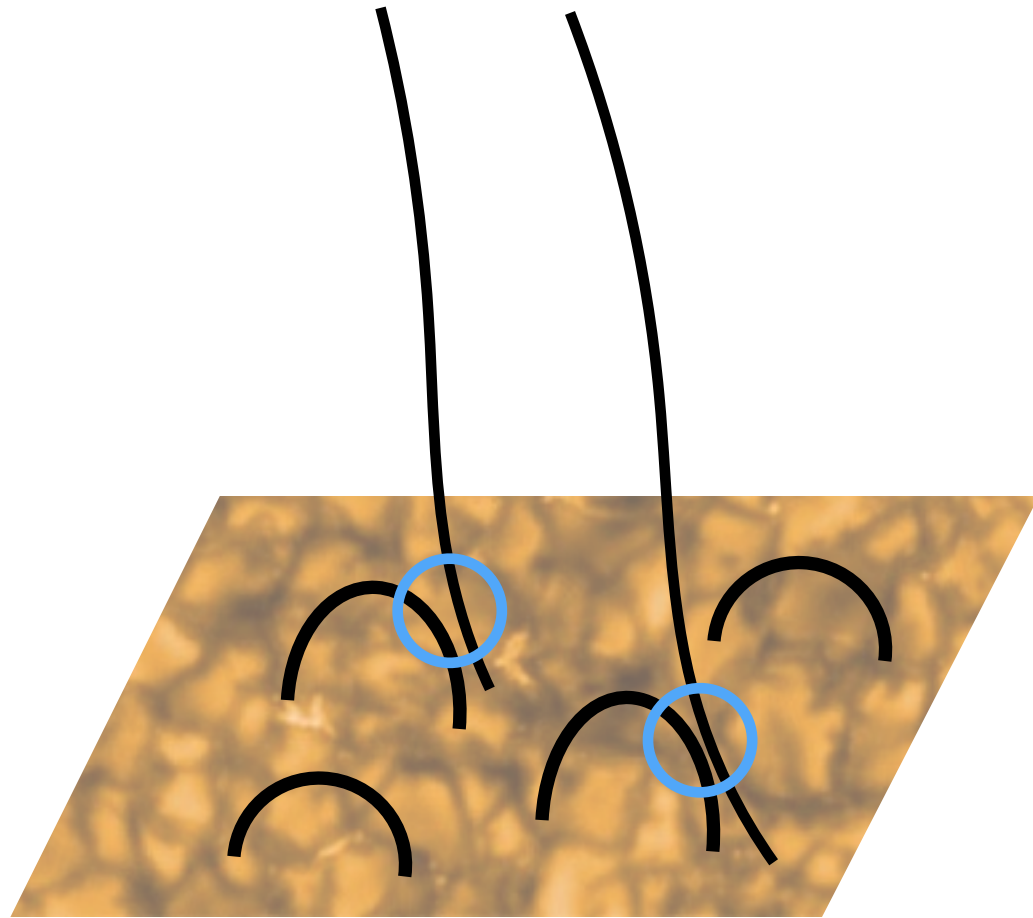
Whip-like Acceleration  
(Yokoyama & Shibata 1996)



x : Reconnection Point  
⇨ : Reconnection Outflow

Takasao+ 2013

# Contribution to Chromospheric and Coronal Heating



Keys for statistical discussion:

- The rate of the energy release by chrom. reconnection
  - the occurrence frequency of chrom. reconnection at various heights
- => Hinode, IRIS, and Solar-C

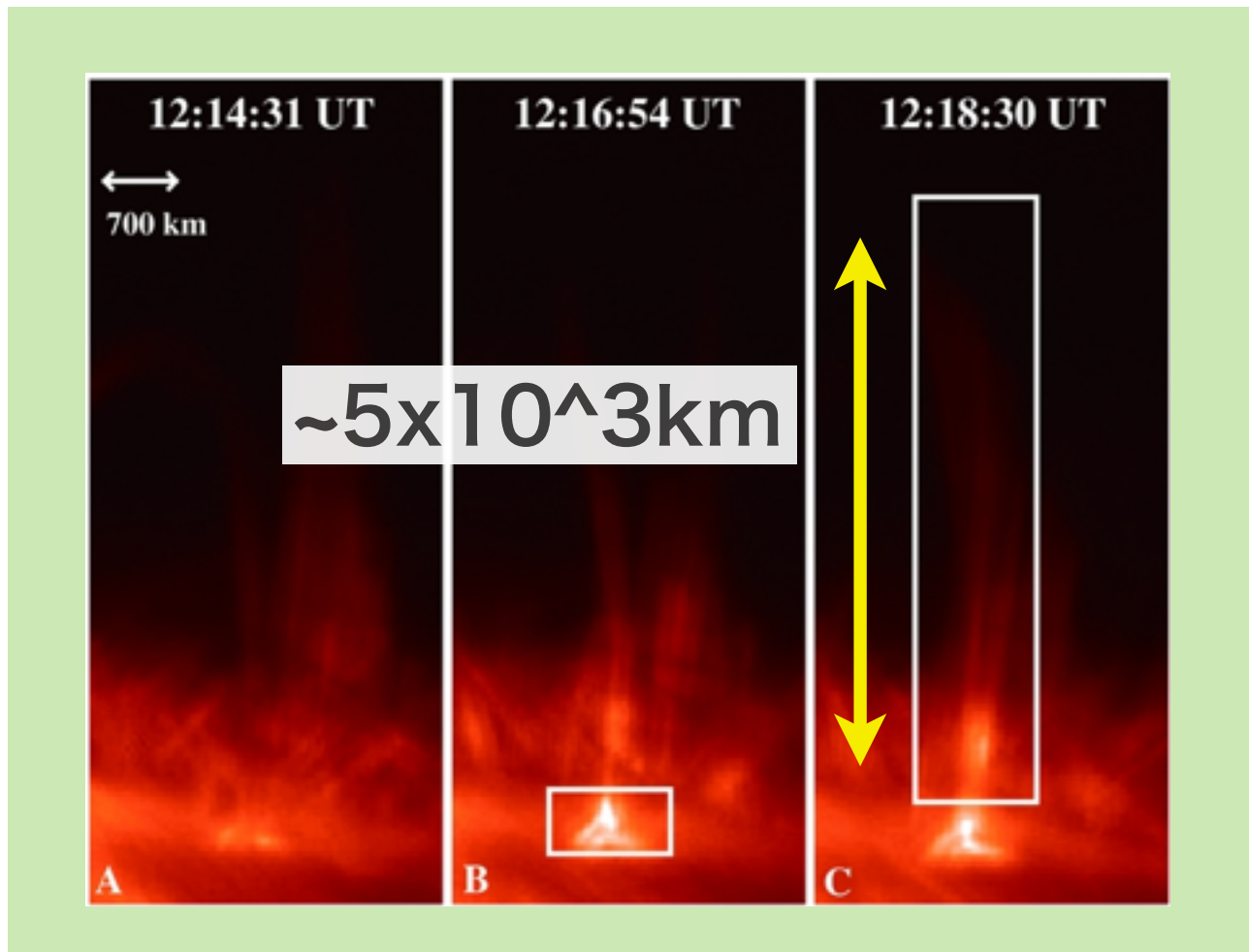
small scale horizontal field  
/granular scale emerging loop  
Centeno+2007, Ishikawa+2008  
Poynting flux  $\sim 10^6 \text{ erg cm}^{-2} \text{ s}^{-1}$   
sufficient to heat the chrom





# Contribution to the Chromospheric Heating

# Can low chromospheric reconnection create tall jets?



(available magnetic energy)  
 ~ (potential energy of jets)

$$\frac{B^2}{8\pi} V \sim \rho g h V$$

V: volume

$$h \sim \frac{B^2}{8\pi} \frac{1}{\rho g}$$

$$= \frac{B^2}{8\pi} \frac{1}{P} \frac{R_g T}{g} = \frac{H_p}{\beta}$$

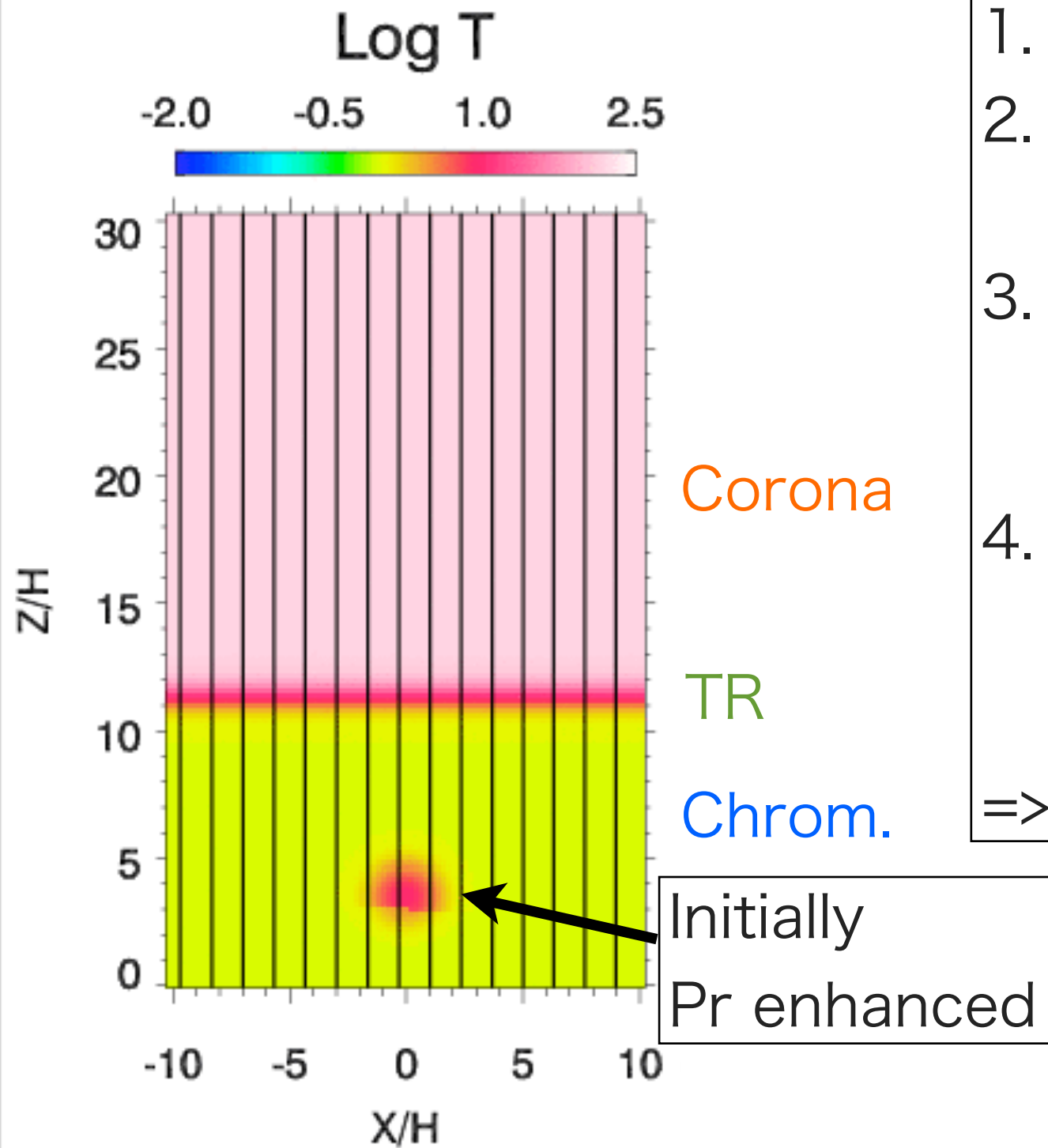
$H_p \sim 200$  km (pressure scale height in the low chrom.)

If  $\beta \sim 1$ ,  $h \sim H_p \sim 200$  km  $\ll 10^3 - 10^4$  km. **Too short!!**

Thus the low chrom. plasma cannot be lifted up to the coronal height by the Lorentz force. So how jets are created?

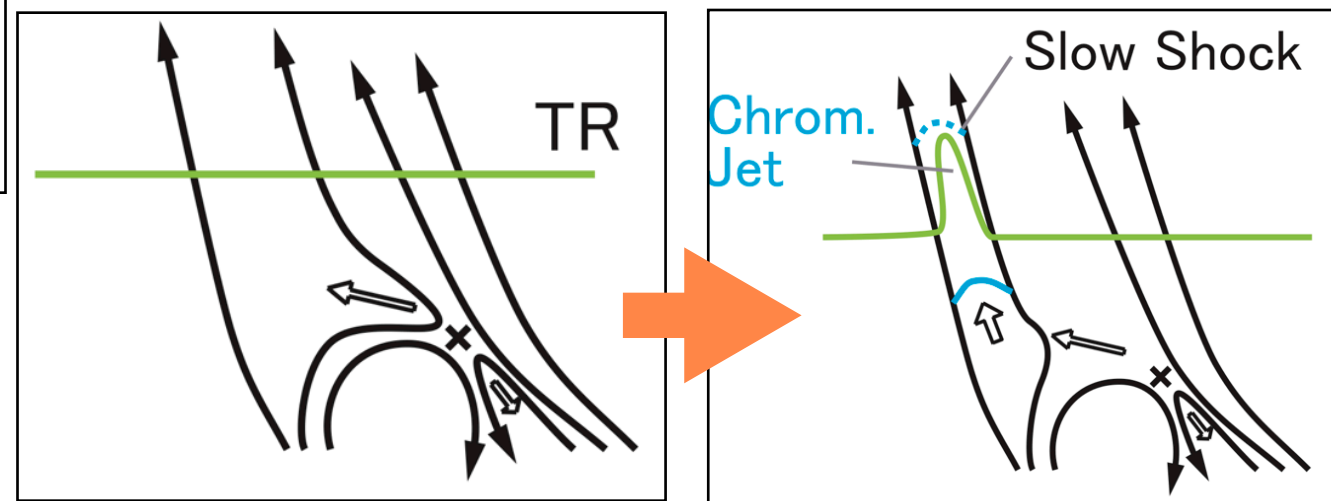


# Energy transport by MHD slow-mode waves/shocks



Solid lines: magnetic field lines

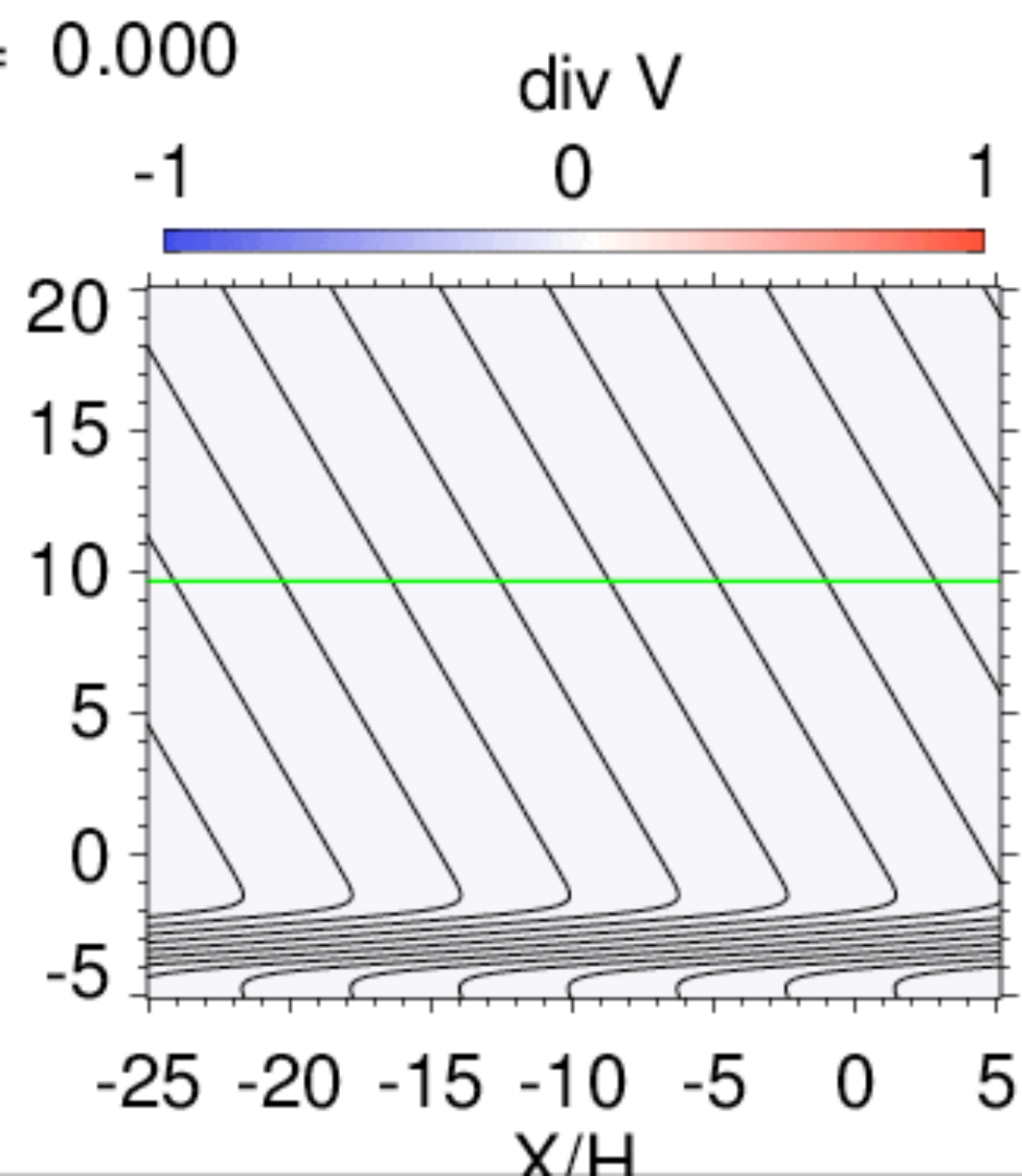
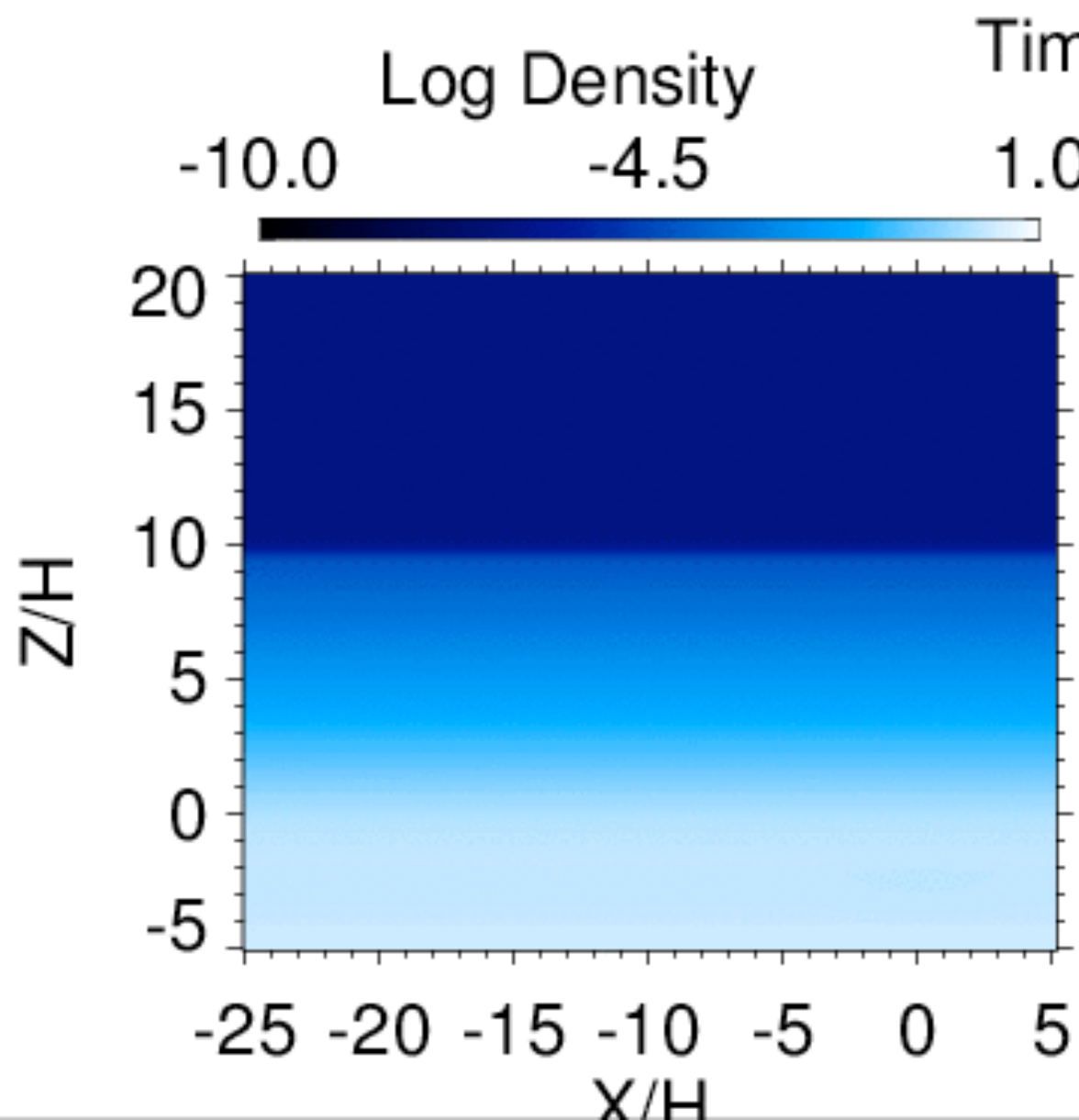
1. Energy release in the low chrom.
2. Slow waves carry the energy along a magnetic field
3. Waves become shocks (e.g. Carlsson and Stein 1997, Centeno + 2009)
4. Only a fraction of the plasma in the upper chrom. (low-density plasma) is accelerated by shocks. => spicules, surges, Ca jets...



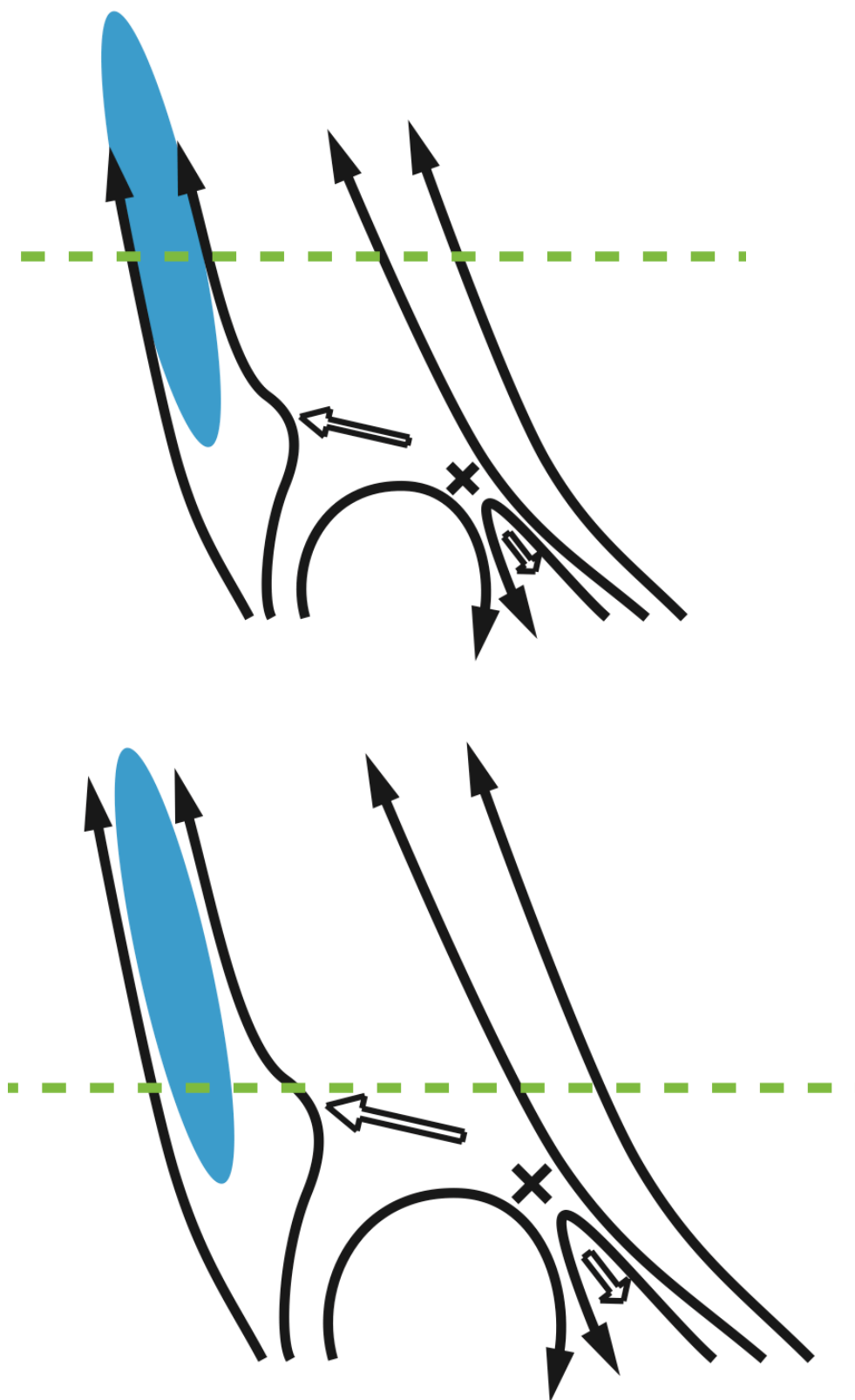
# Density and div V distributions

divergence  
of the velocity field

Blue : compressed (~shocks)  
Red : expanded



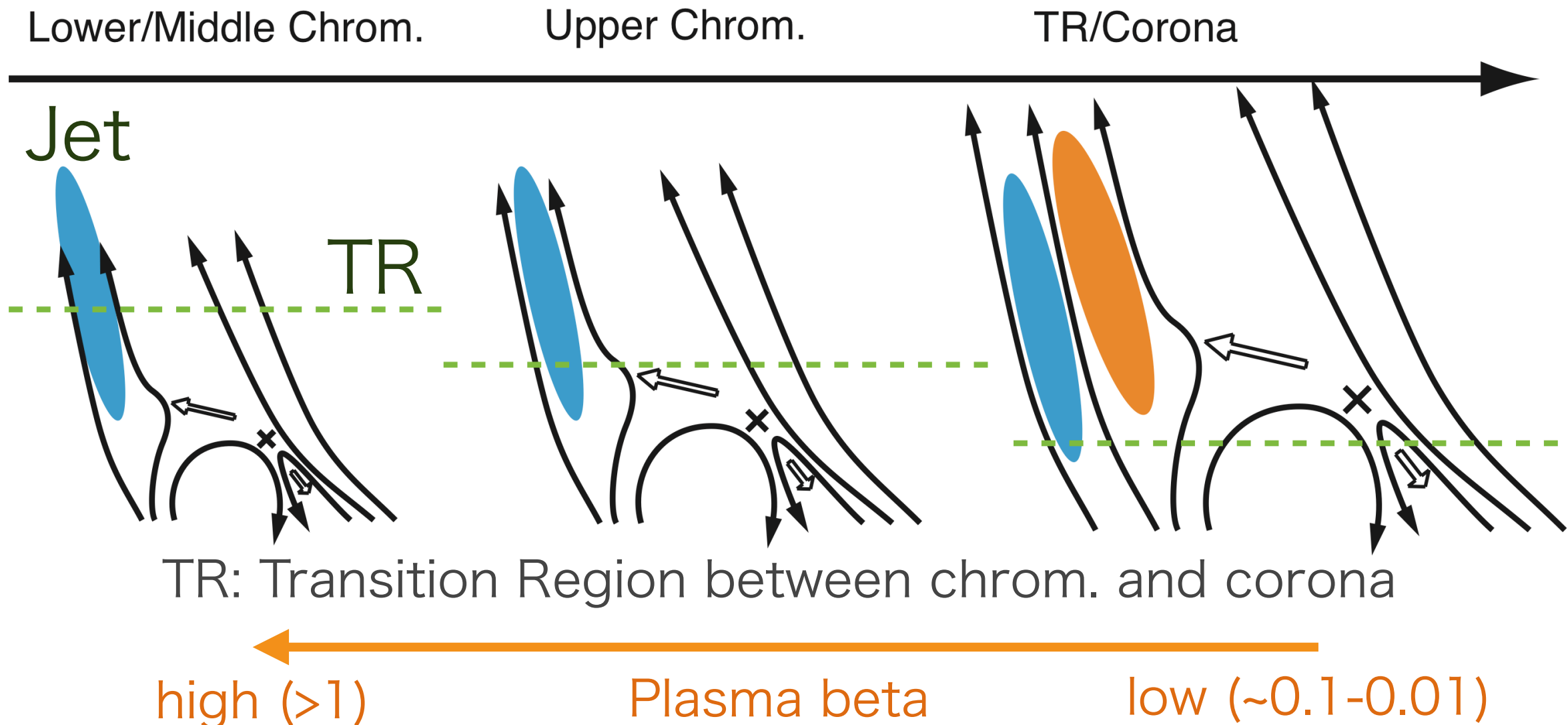
# Magnetic reconnection can take place at various heights



1. Energy release in the low chrom.
2. Slow waves carry the energy along a magnetic field
3. Waves become shocks (e.g. Shibata + 1982, Carlsson and Stein 1997, Centeno + 2009)
4. Only a fraction of the plasma in the upper chrom. (low-density plasma) is accelerated by shocks.  
=> spicules, surges, Ca jets...

# Magnetic Reconnection can Take Place at Various Heights

## Height of Reconnection point



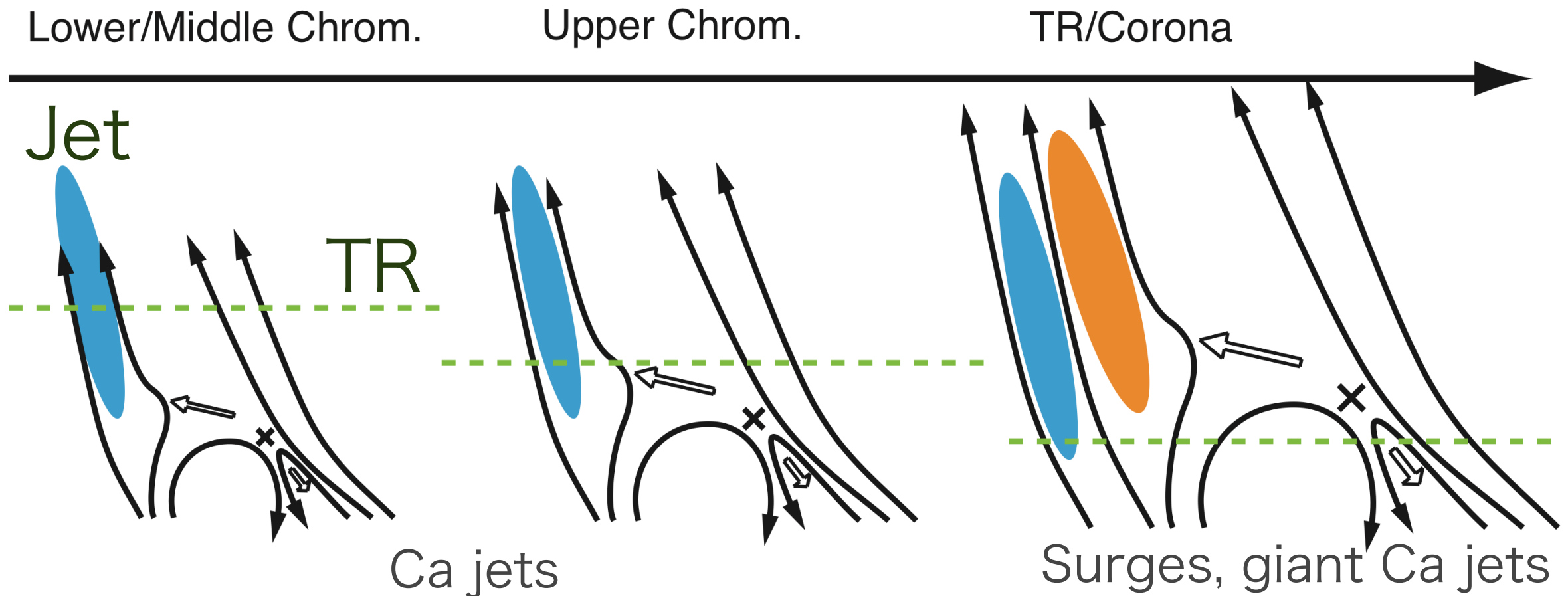
Their apparent structures are quite similar.

=> Their magnetic structures are similar.

=> Same acceleration scenario? NO!

# Magnetic Reconnection can Take Place at Various Heights

## Height of Reconnection point

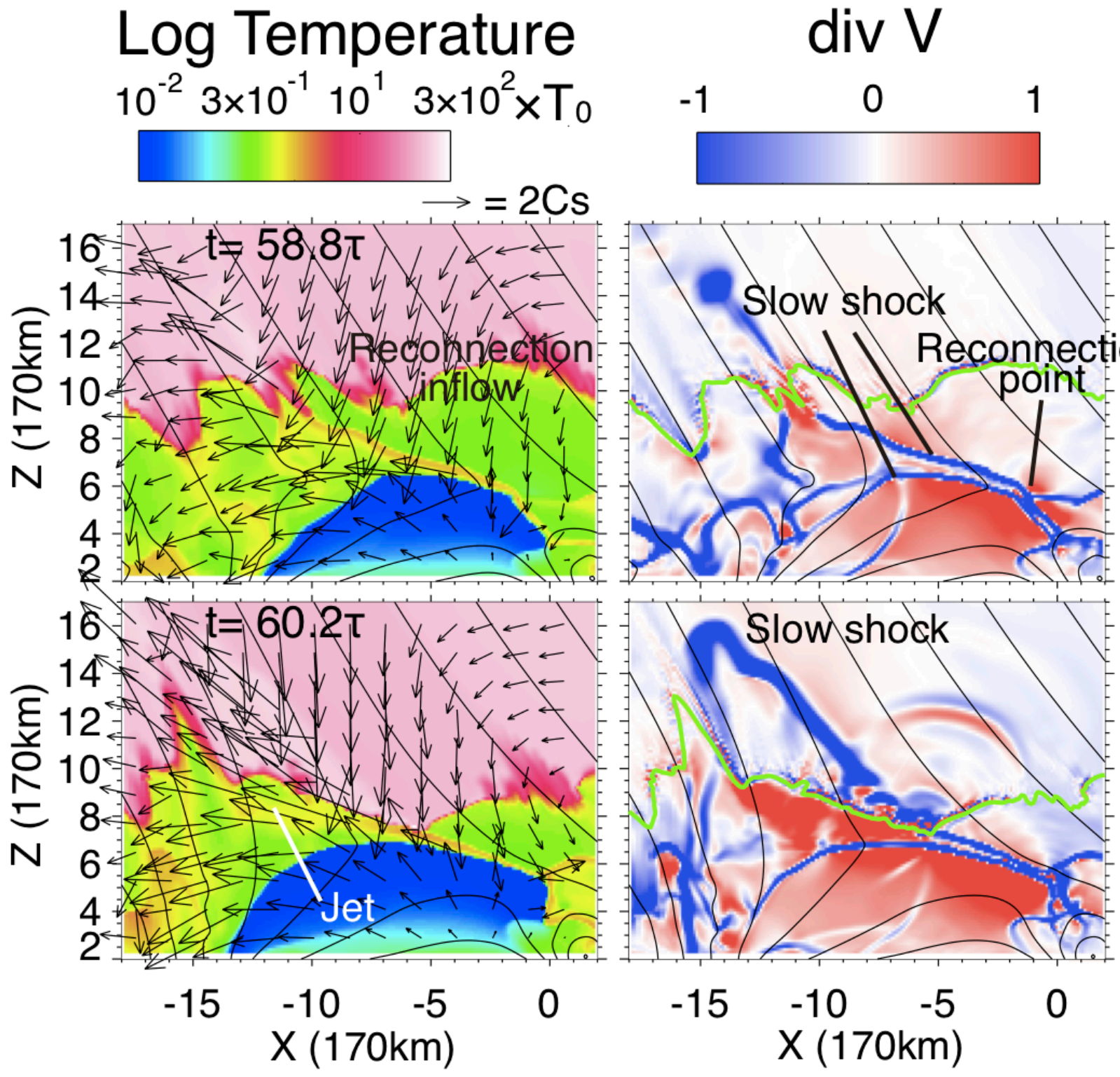


TR: Transition Region  
between chrom. and corona

Sling-shot  
or **whip-like acceleration**  
e.g. Yokoyama&Shibata 1996,  
Moreno-Inertis+ 2013

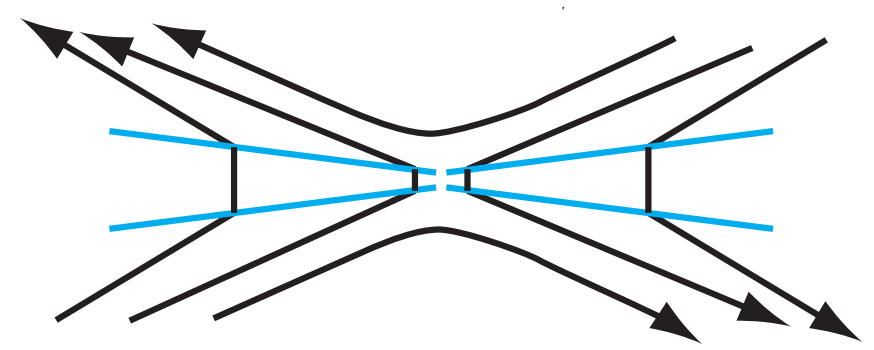


# Petschek Slow Shock Contributes to Acceleration of Jet



Before

Shock crosses TR



After

Shock crosses TR

A Petschek shock crosses TR, and then accelerates the plasma.



# Chromospheric Jets

- ▶ Spicules      small
- ▶ **Ca jets**      ↓
- ▶ Surges      large

Ca jets

time scale:

~a few-10 min.

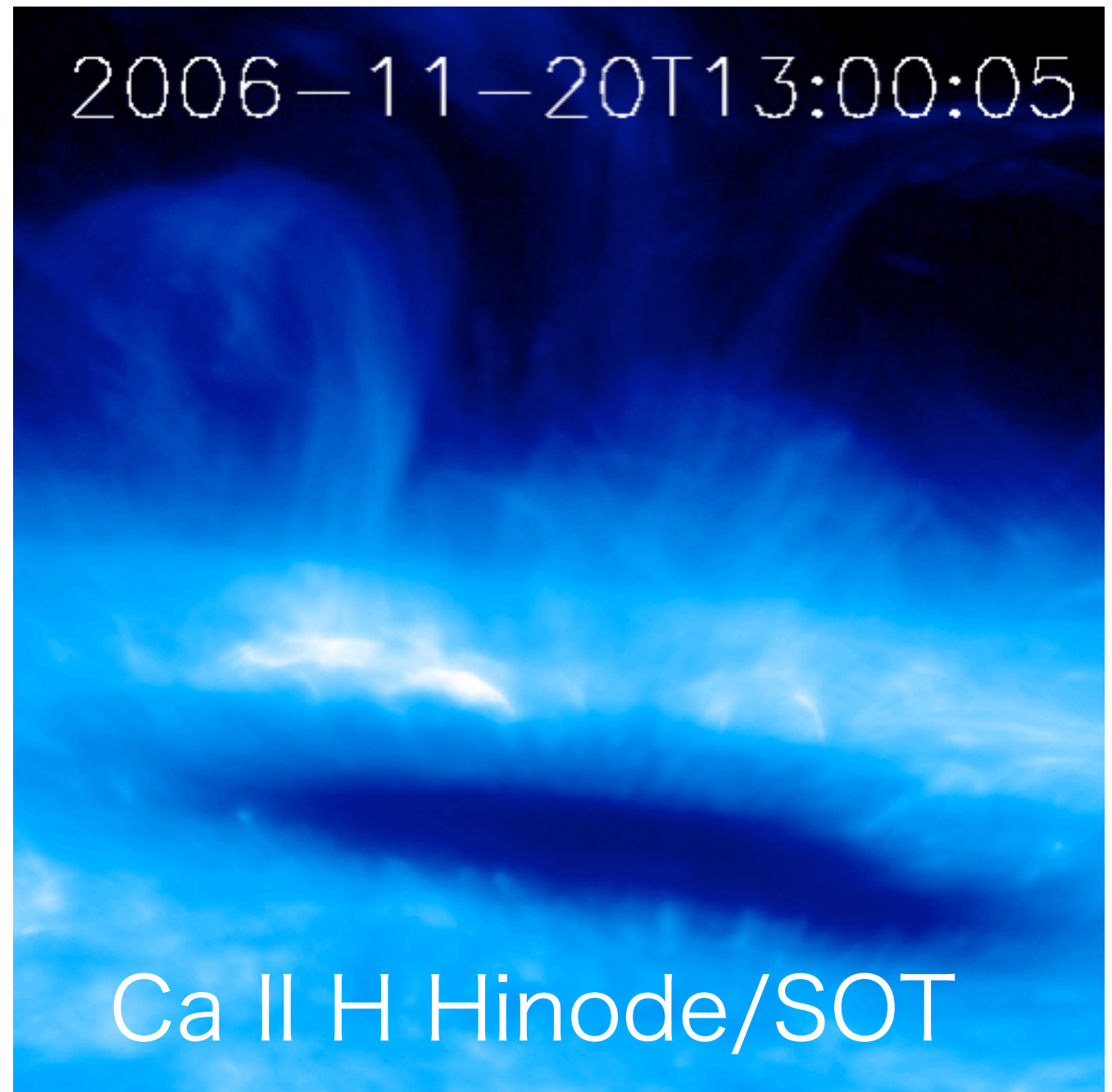
length scale:

$10^3 - 10^4$  km

velocity:

5 - 20 km/s

(Nishizuka+2011)



e.g. Sterling + 1993, Shibata + 2007  
De Pontieu + 2007, Nishizuka + 2008  
Yang + 2013, Kayshap + 2013

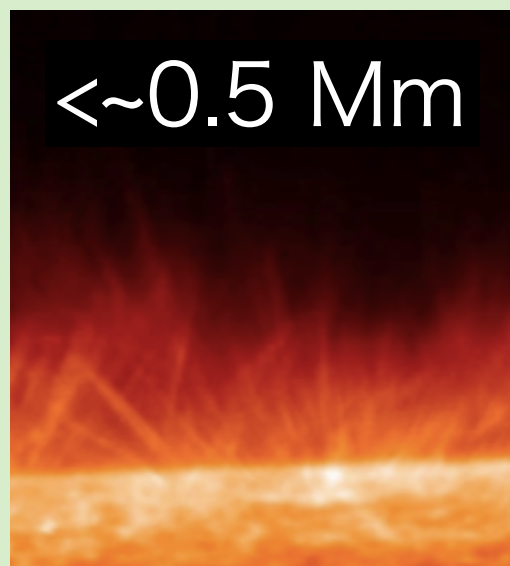


# Classification by Height of Drivers of Jets

## Height of Drivers of Jets

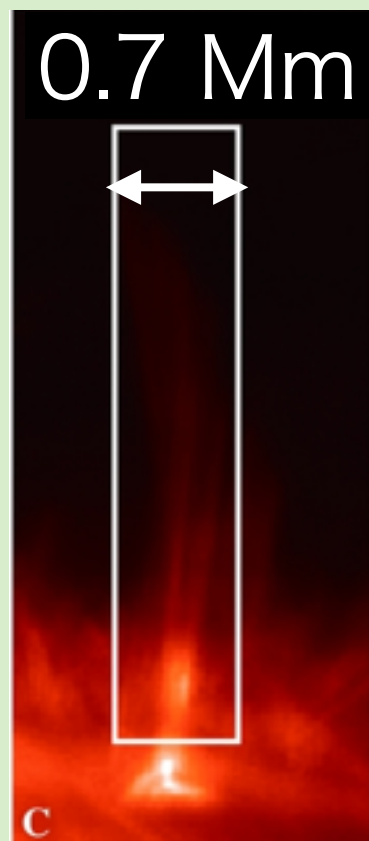
Photo./Low chrom.      Upper chrom.      TR/corona

Spicules

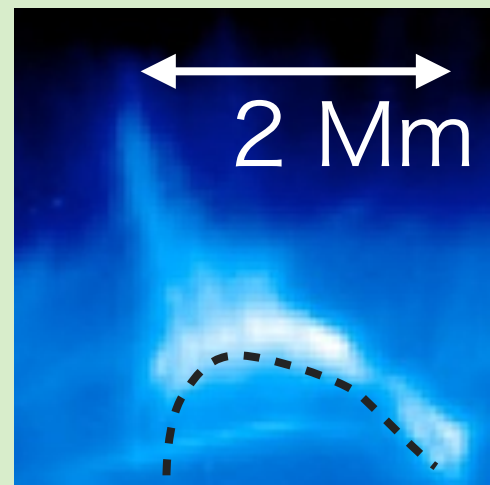


foot-points  
unresolved...

Ca jets



Singh+2012

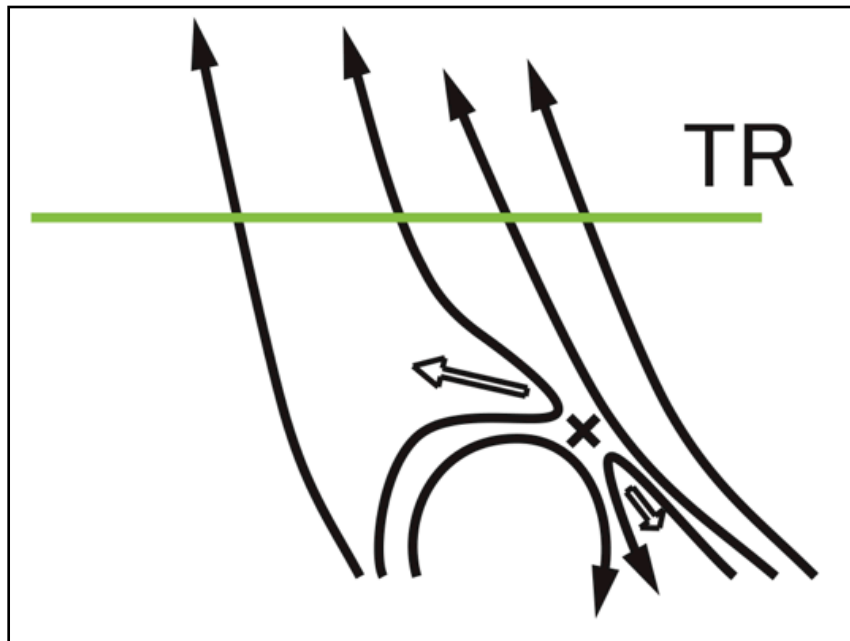


Giant Ca jets / Surges



loop- or cusp-shaped foot-points  
=> probably emerging flux

# Low Atmospheric Reconnection and Shock



Height of jets driven by the Lorentz force  
(sling-shot / whip-like acceleration)

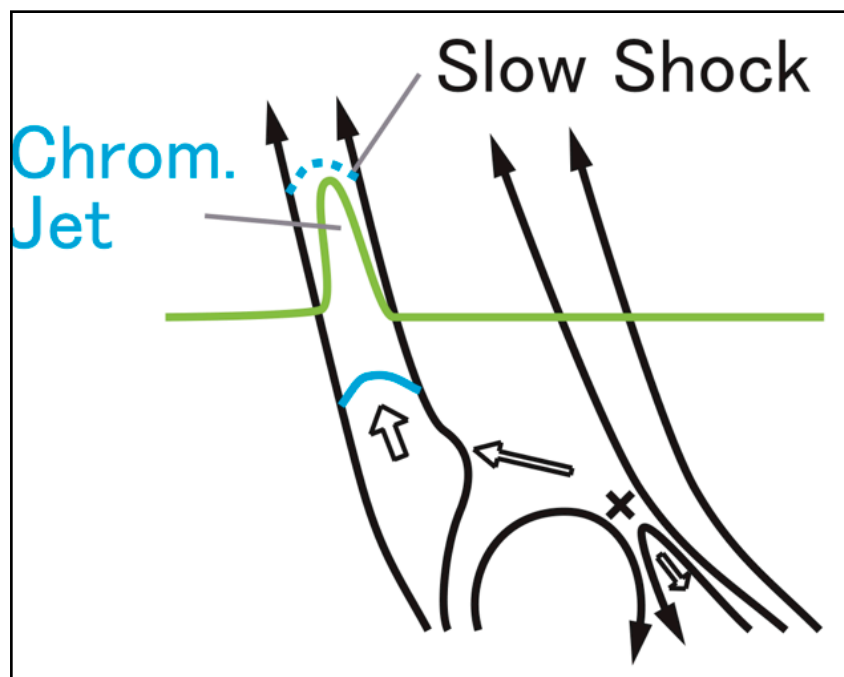
$$h_{jet} \sim H_p / \beta \sim 150 / \beta \text{ km}$$

$$h_{jet,obs} \sim 1 - 4 \times 10^3 \text{ km}$$

OK for low- $\beta$  plasma (corona)

No for high- $\beta$  plasma (low chrom.)

=> Shock acceleration



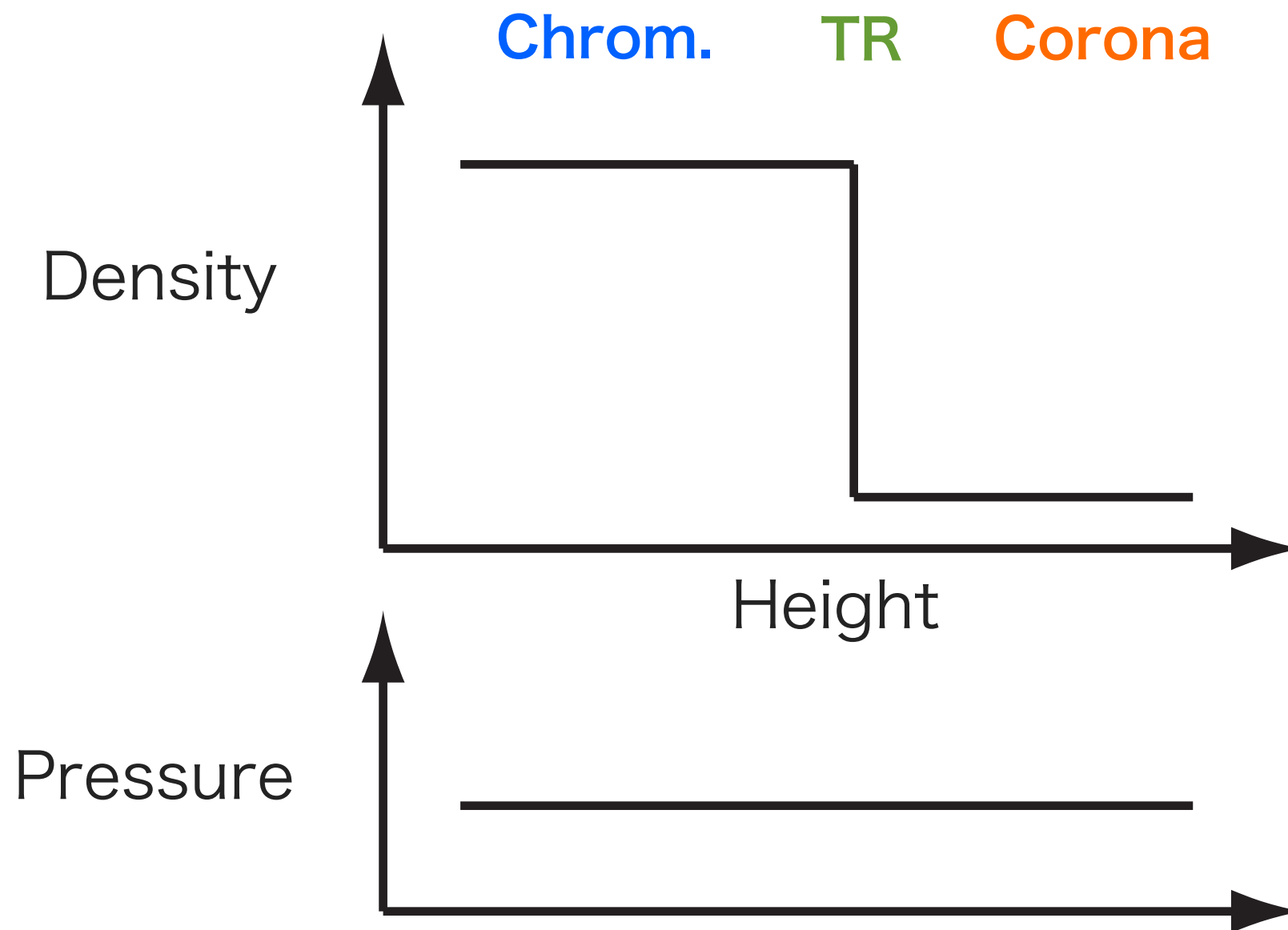
1. Energy release in the low chrom.
2. Slow waves/shocks carry the energy along a magnetic field
3. Only a fraction of the plasma in the upper chrom. (low-density plasma) is accelerated by shocks.

(e.g. Shibata+1982, 2007,  
Hegglund+2007)

# How Shocks Create Jets?

## Transition Region - Shock Interaction

When a shock passes through TR, TR is launched to become a chrom. jet.

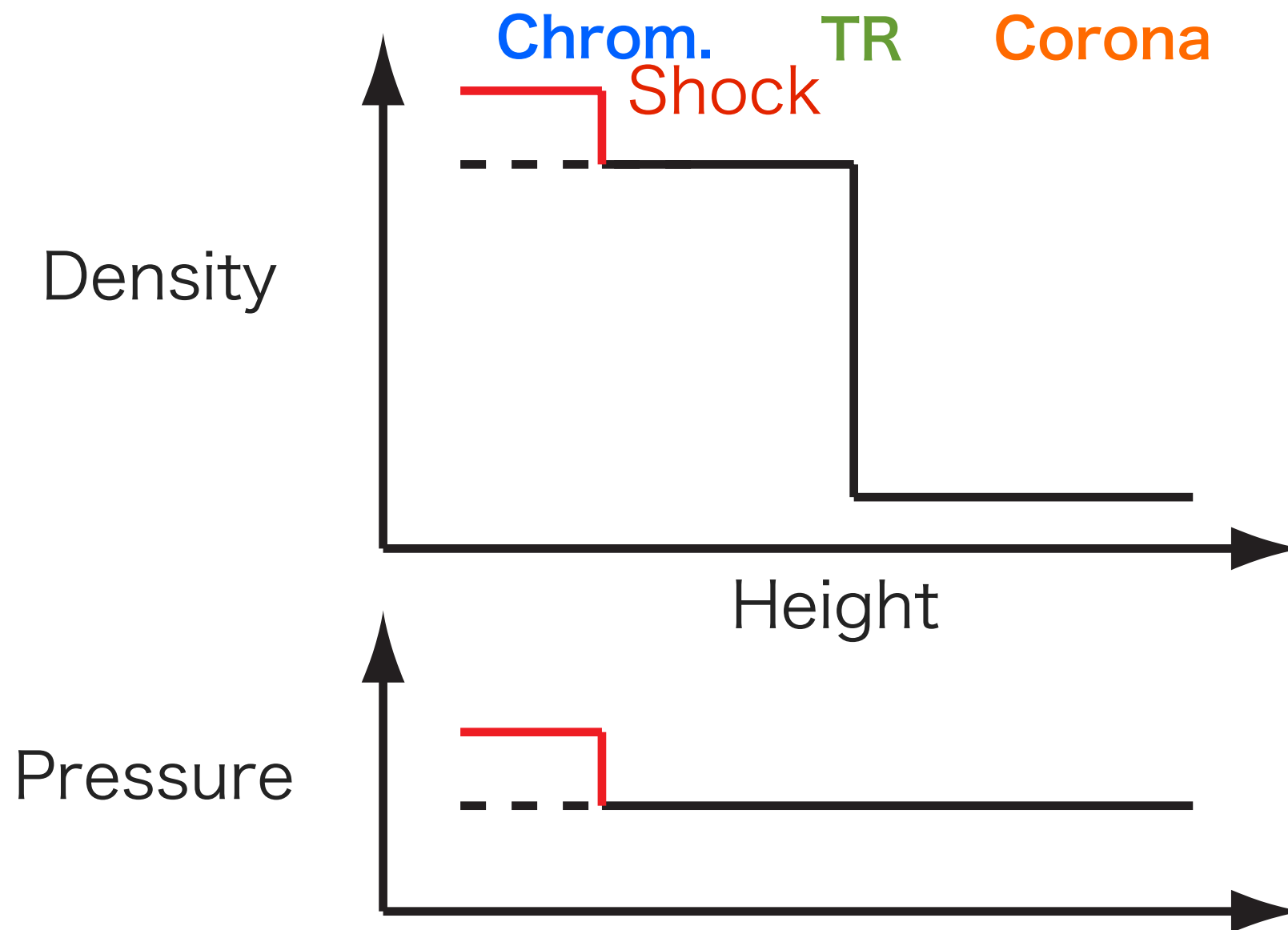


Note: Here the stratification is neglected

# How Shocks Create Jets?

## Transition Region - Shock Interaction

When a shock passes through TR, TR is launched to become a chrom. jet.



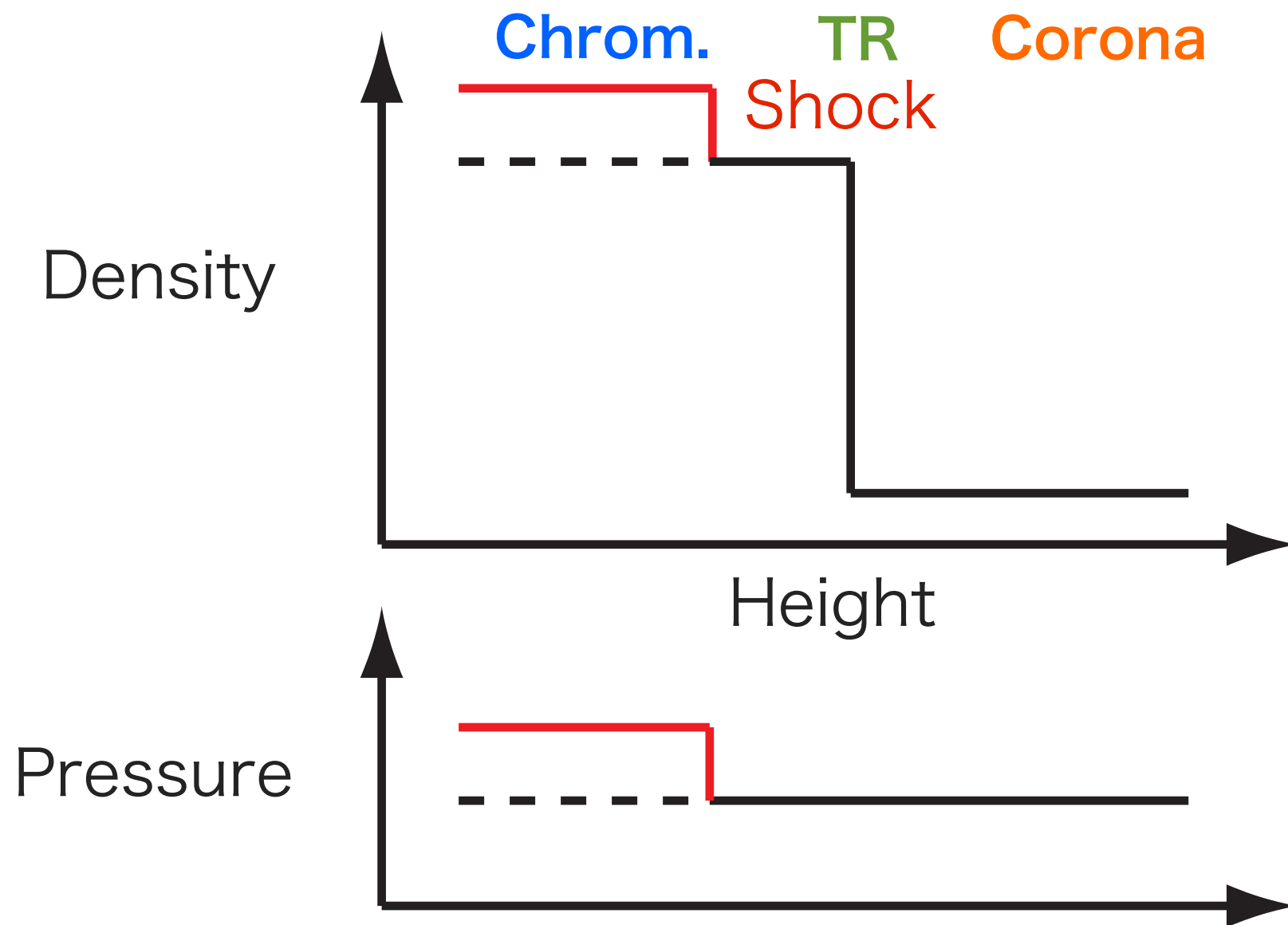
Note: Here the stratification is neglected



# How Shocks Create Jets?

## Transition Region - Shock Interaction

When a shock passes through TR, TR is launched to become a chrom. jet.

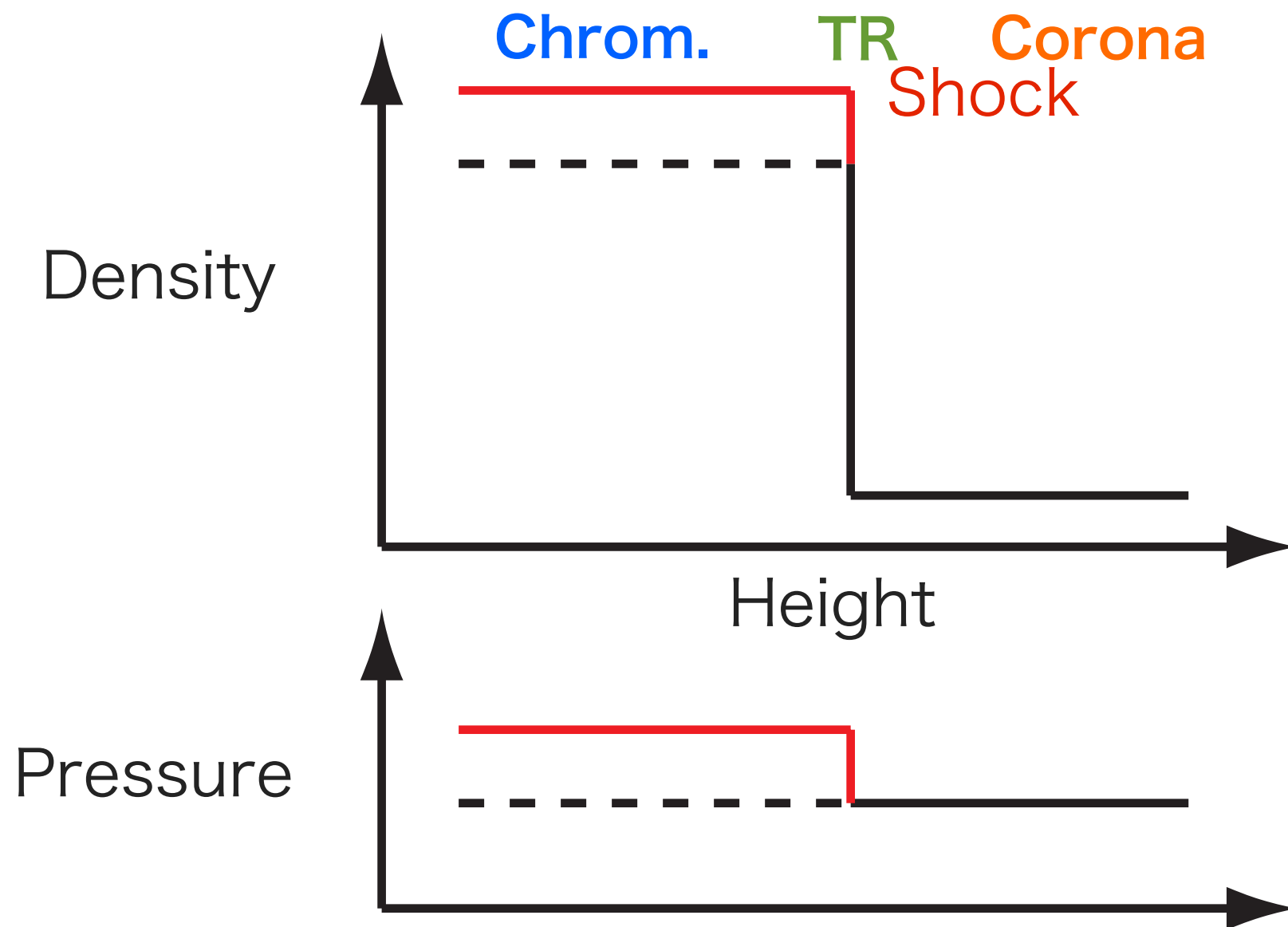


Note: Here the stratification is neglected

# How Shocks Create Jets?

## Transition Region - Shock Interaction

When a shock passes through TR, TR is launched to become a chrom. jet.

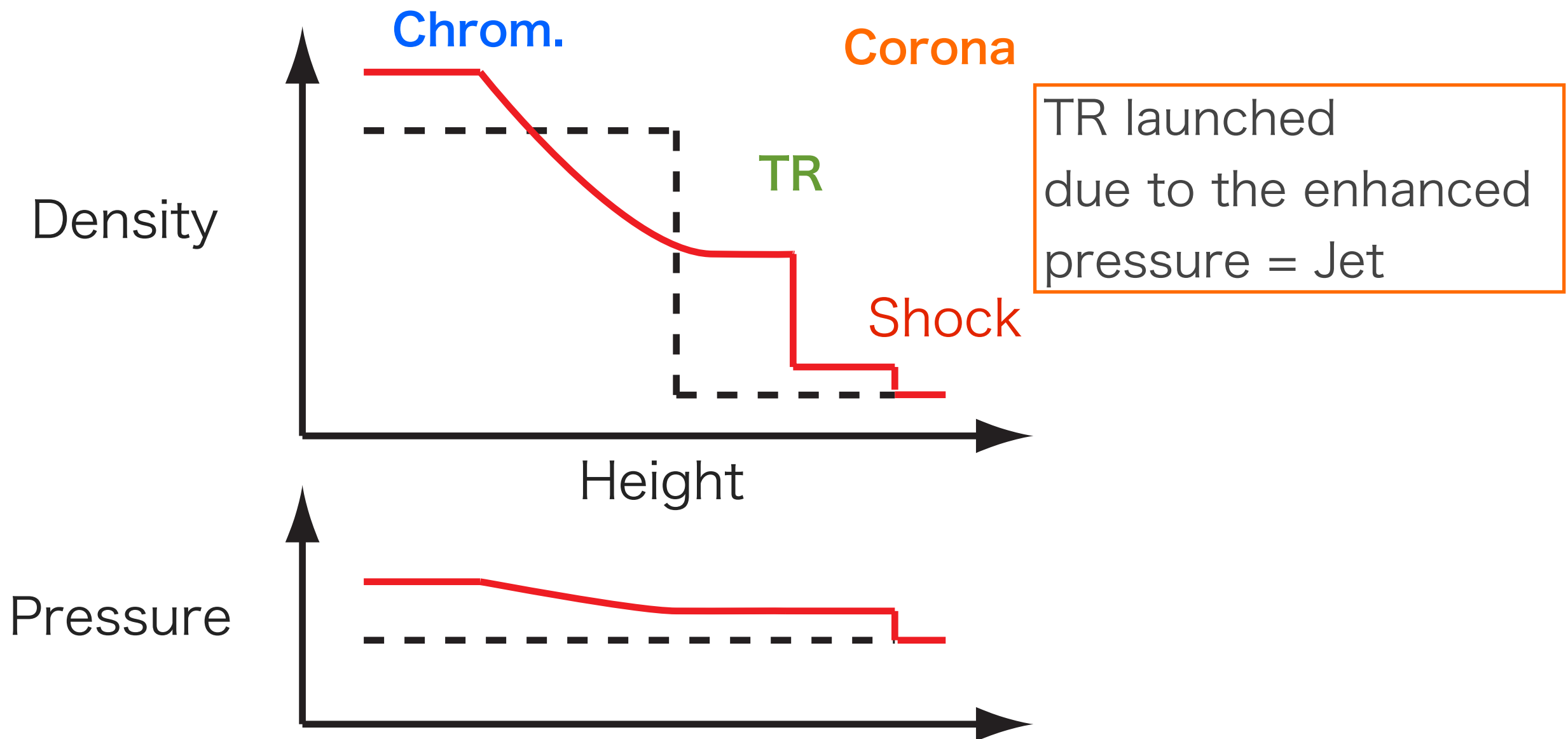


Note: Here the stratification is neglected

# How Shocks Create Jets?

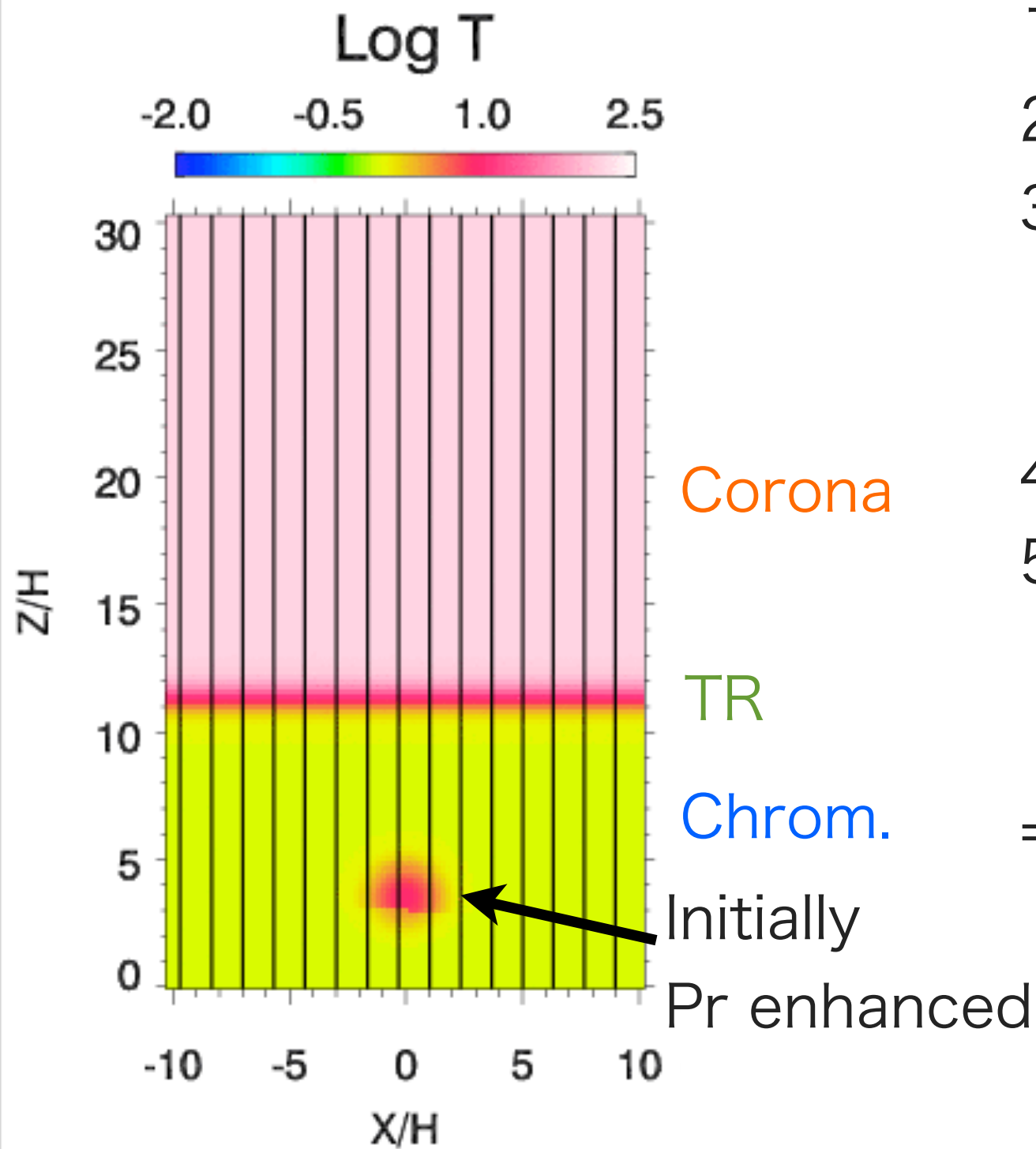
## Transition Region - Shock Interaction

When a shock passes through TR, TR is launched to become a chrom. jet.



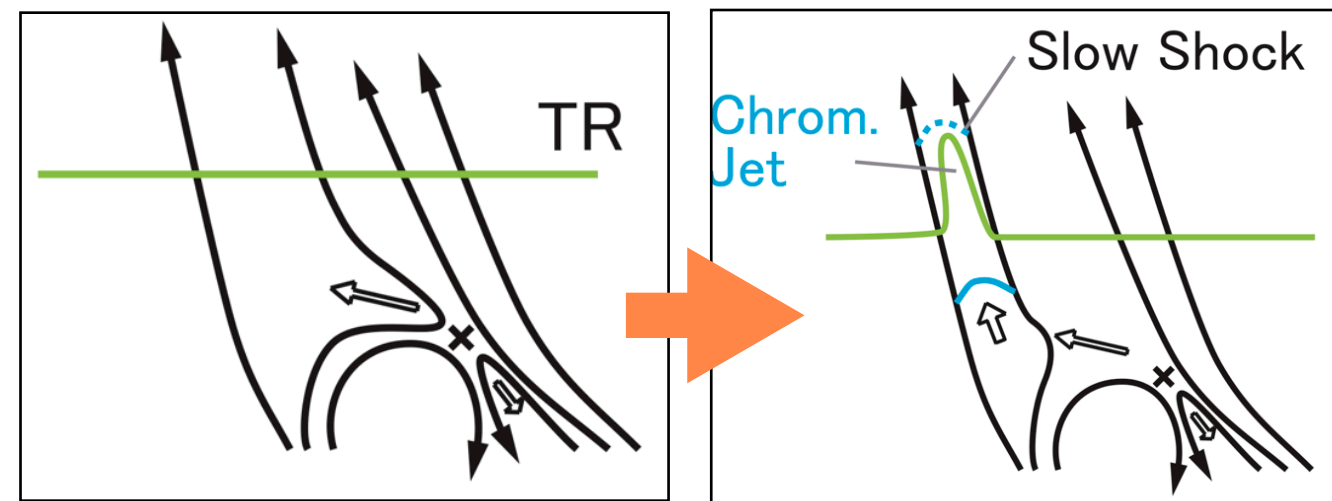
Note: Here the stratification is neglected

# An Example of Jet Created by a Shock (=>Jets due to Reconnection in a Low Chrom.)



Solid lines: magnetic field lines

1. Energy release in the low chrom.
  2. Waves carry the energy
  3. Due to the stratification, the amplitude of the waves drastically increases
  4. Finally waves become shocks
  5. Only a fraction of the plasma in the upper chrom. (low-density plasma) is accelerated by shocks.
- => spicules, surges, small Ca jets...



$$\frac{\partial \rho}{\partial t} + (\mathbf{v} \cdot \nabla) \rho = -\rho \nabla \cdot \mathbf{v}$$

$$\frac{\partial v}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{1}{\rho} \nabla p + \frac{1}{4\pi\rho} (\nabla \times \mathbf{B}) \times \mathbf{B} + \mathbf{g}$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B} - \eta \mathbf{J})$$

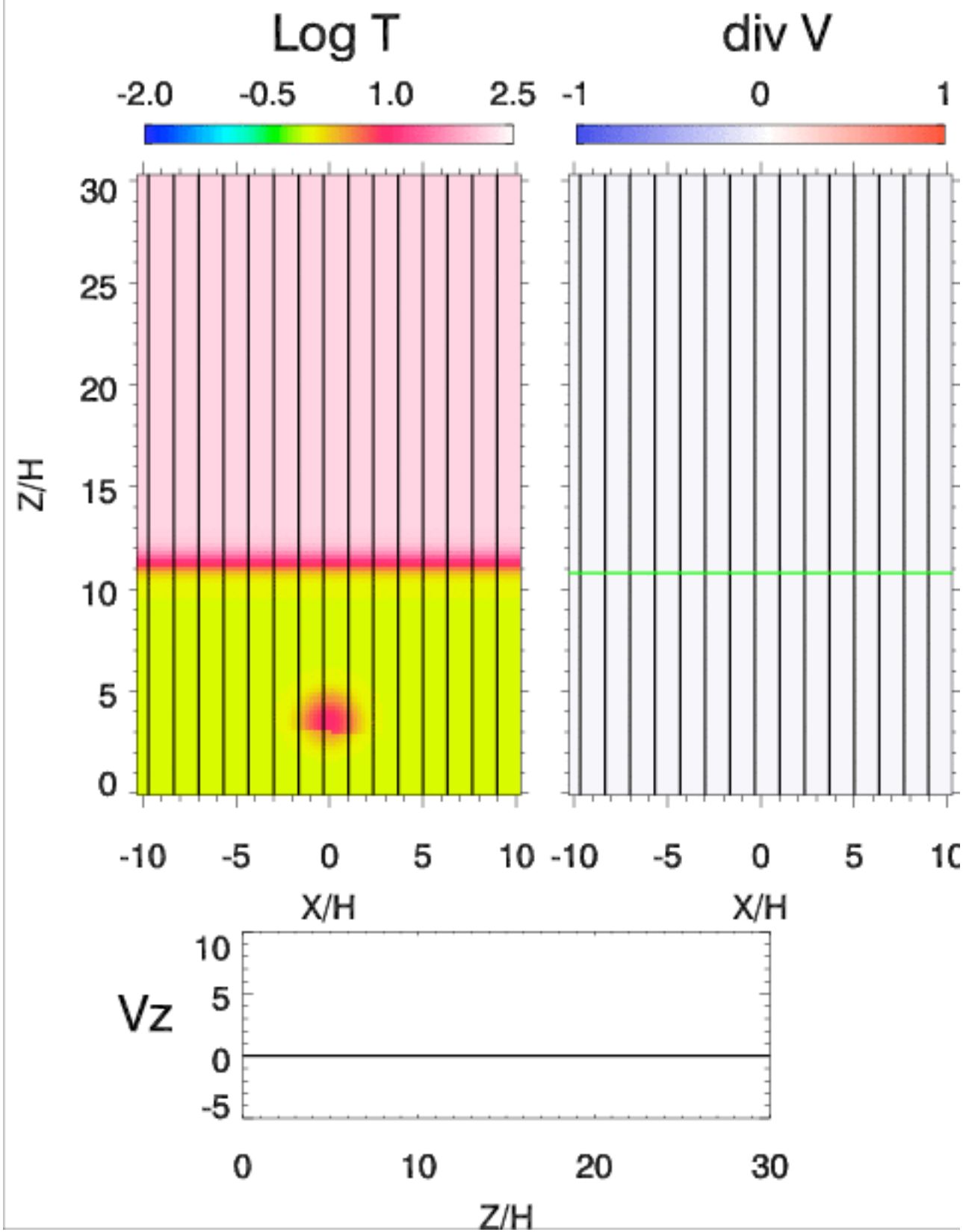
$$\mathbf{J} = \frac{1}{4\pi} \nabla \times \mathbf{B}$$

$$\frac{\partial T}{\partial t} + (\mathbf{v} \cdot \nabla) T = -(\gamma - 1) T \nabla \cdot \mathbf{v} - \frac{T}{\tau_{cooling}}$$

$$p = \frac{k_B}{m} \rho T$$

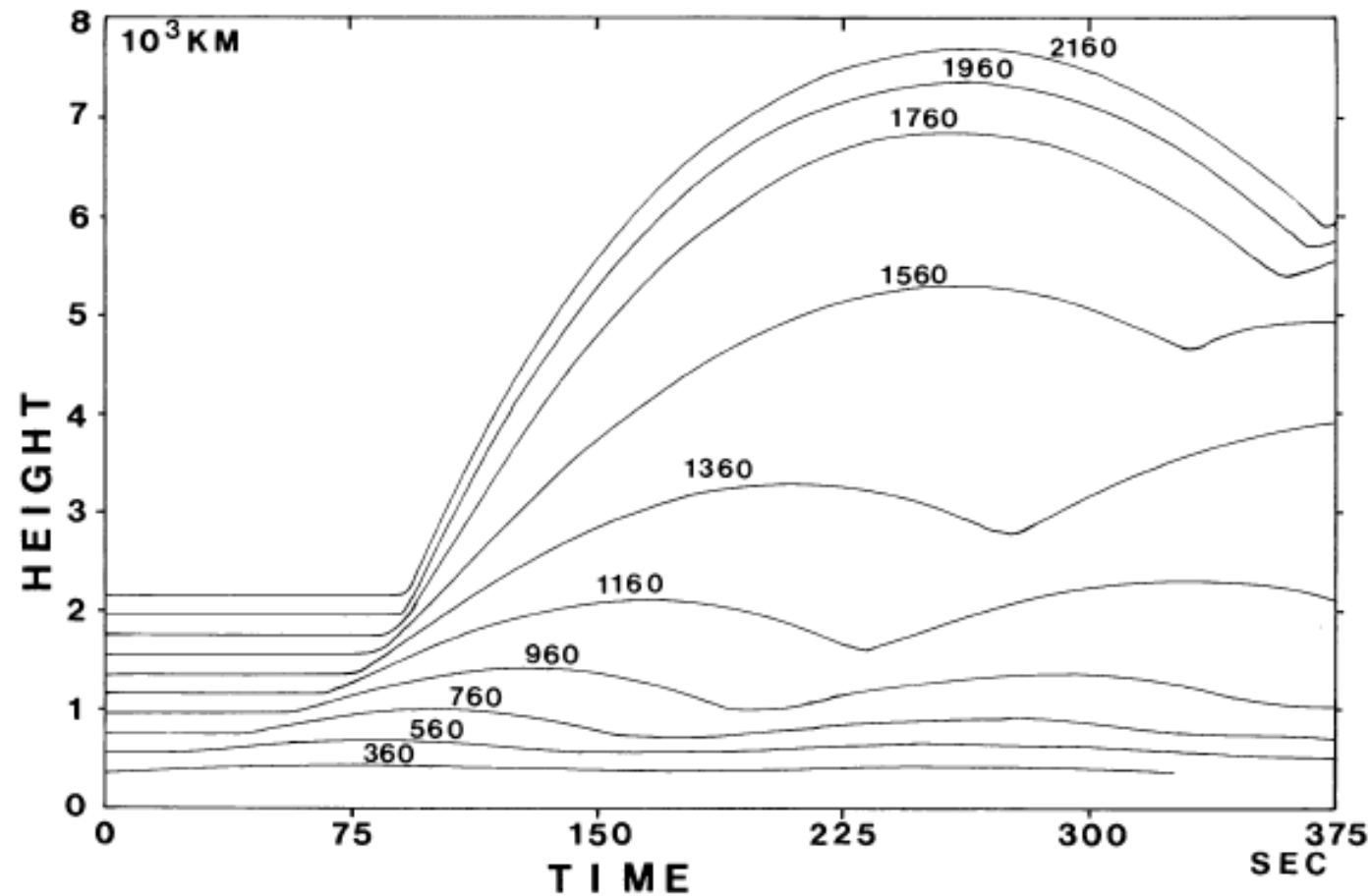
$$\eta = \begin{cases} 0 & \text{for } v_d < v_c \\ \alpha(v_d/v_c - 1)^2 & \text{for } v_d \geq v_c \end{cases} \quad \tau_{cooling}(z) = \begin{cases} \infty & (z < 0) \\ \frac{\tau_{c1} - \tau_{c0}}{z_{tr}} z + \tau_{c0} & (0 \leq z \leq z_{tr}) \\ \tau_{c1} & (z > z_{tr}), \end{cases}$$

# Point Explosion





# Shock acceleration

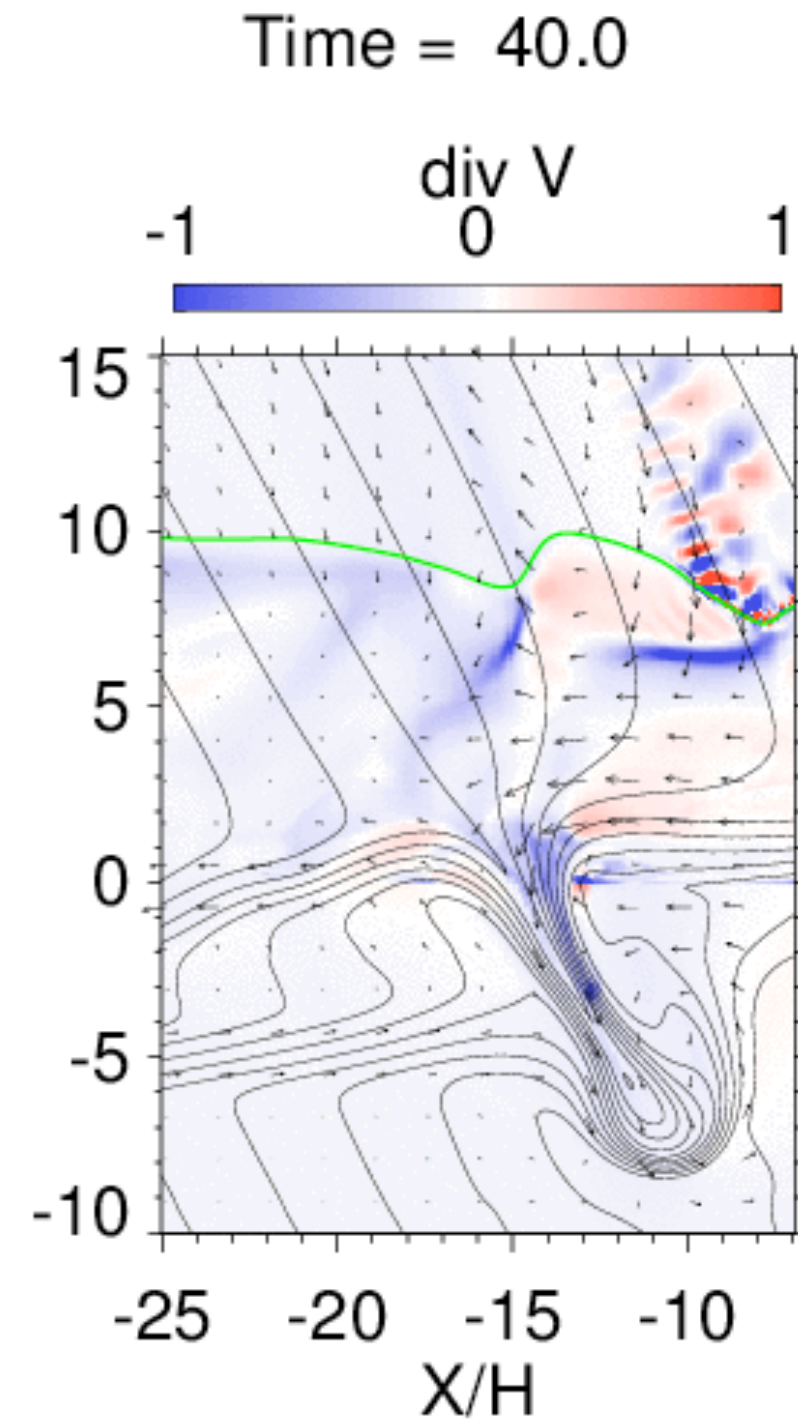
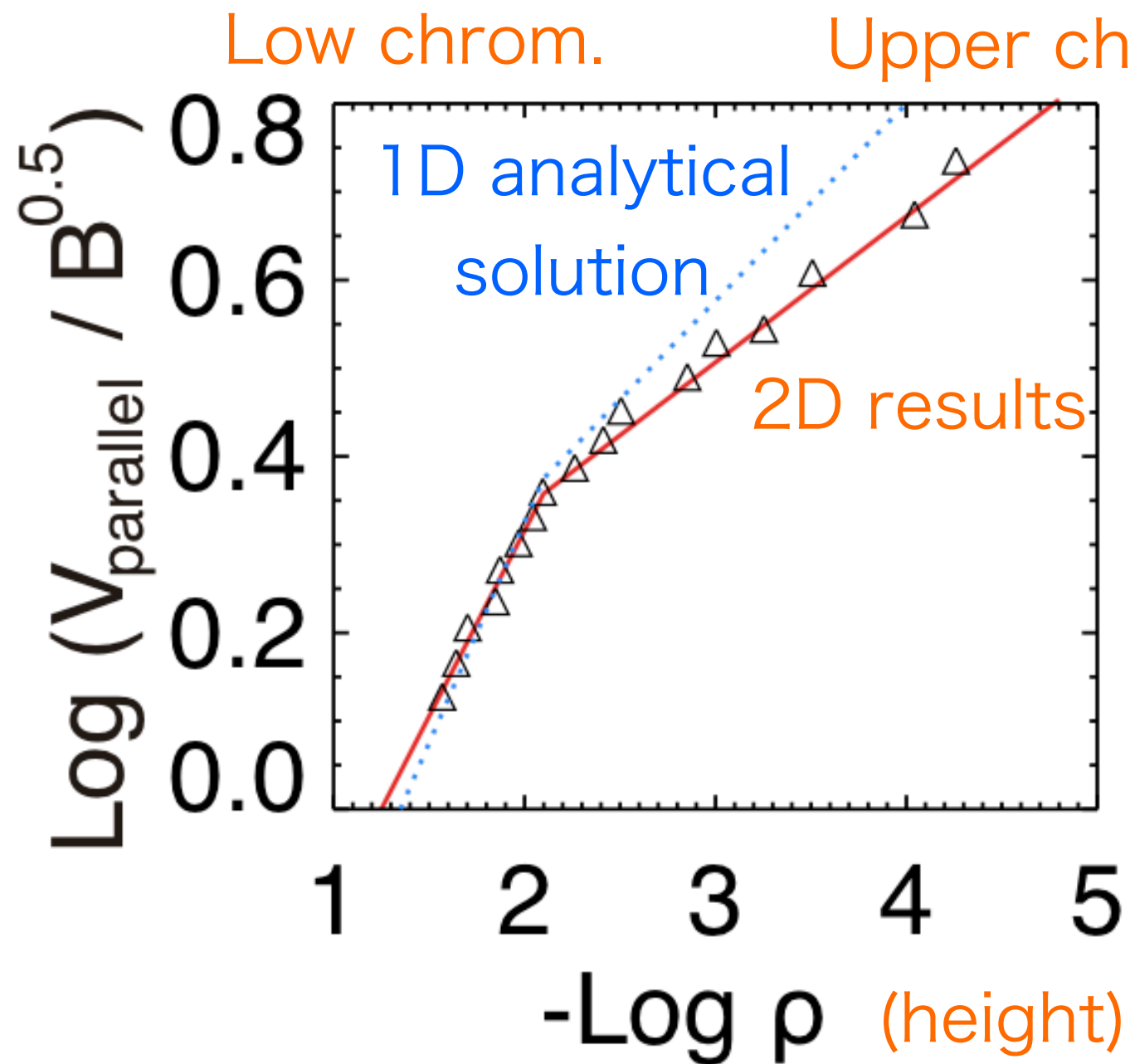


Trajectories  
of Lagrange particles

Fig. 2(a). The trajectories of the fluid elements which initially exist at the height of 360, 560, 760, 960, 1160, 1360, 1560, 1760, 1960, and 2160 km. Numerals beside each curve represent the initial heights of the fluid elements.

Suematsu et al. 1982

# Growth of the Amplitude of a Shock



ambipolar diffusion (current sheet thinning)

$$\frac{\partial B}{\partial t} = \nabla \times \left[ V_n \times B - \frac{J \times B}{en_e} + \frac{(J \times B) \times B}{cV_{ni}\rho_n} - \eta J \right]$$

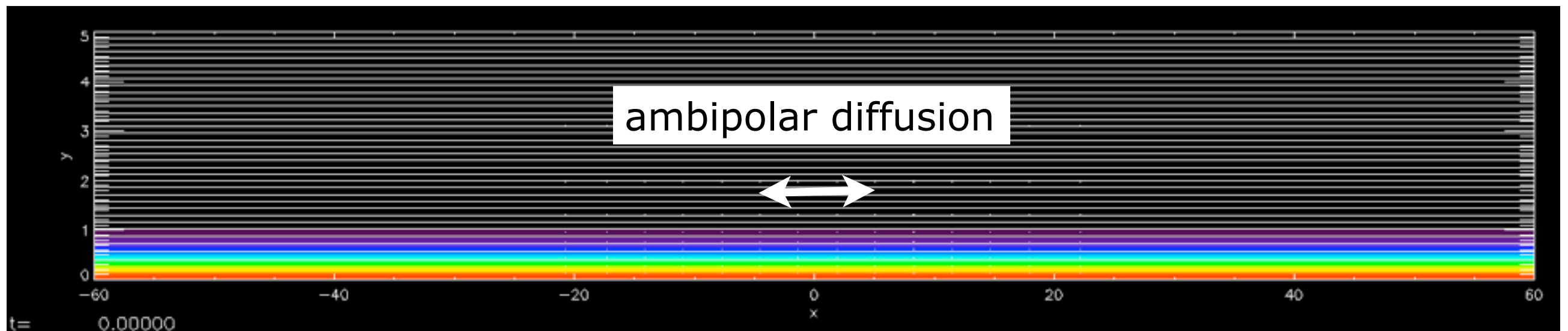
localized distribution of neutrals



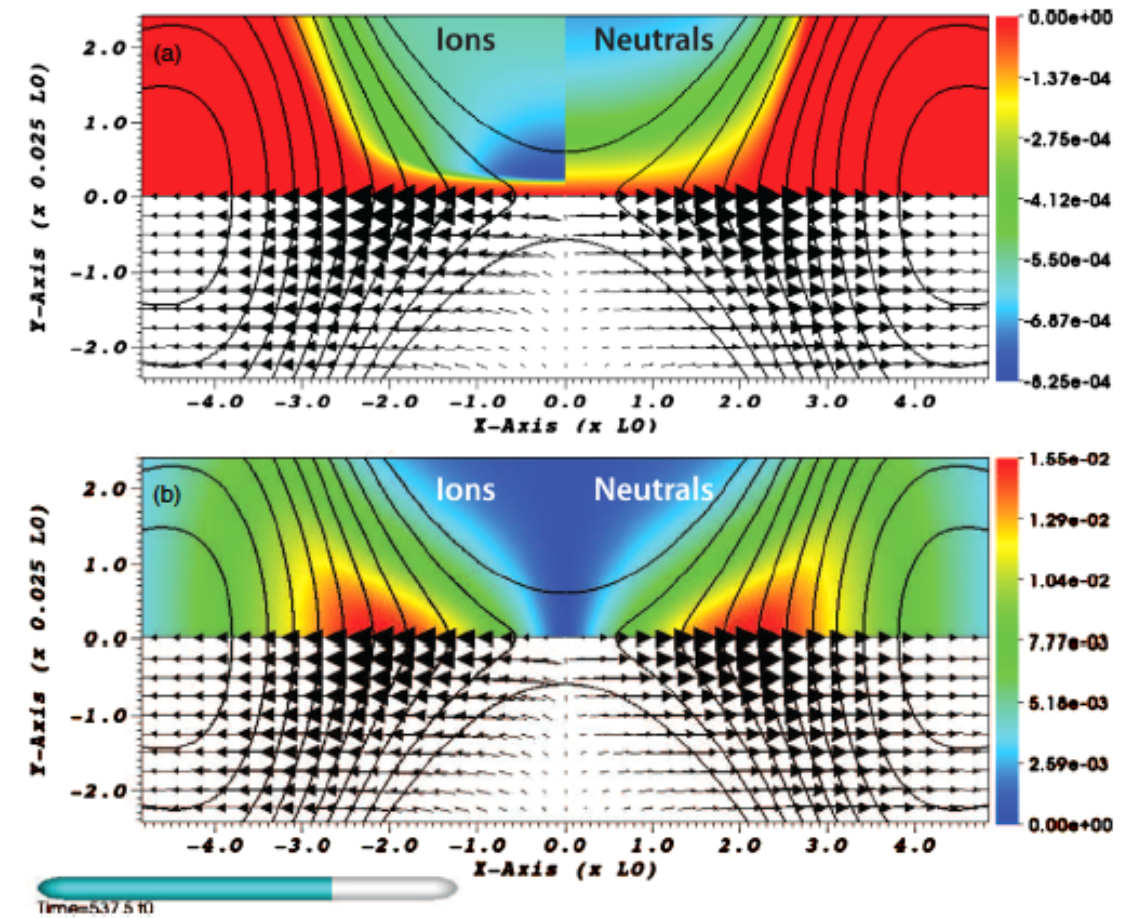
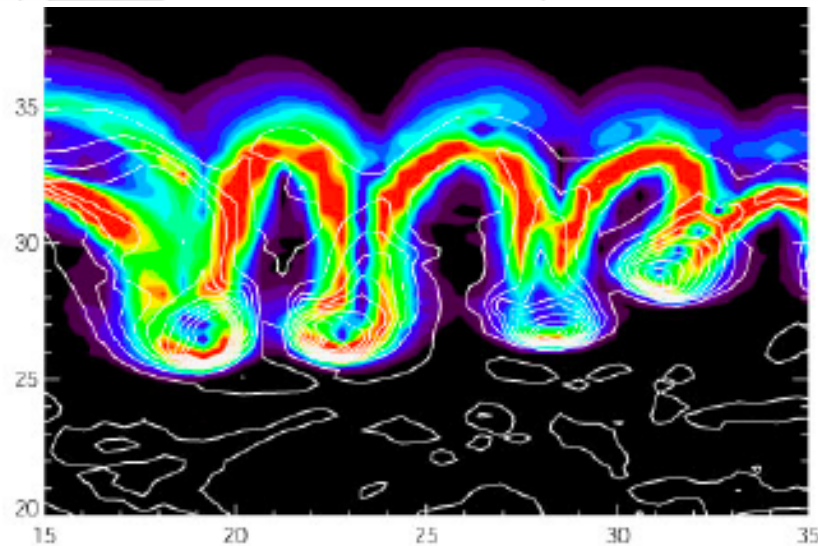
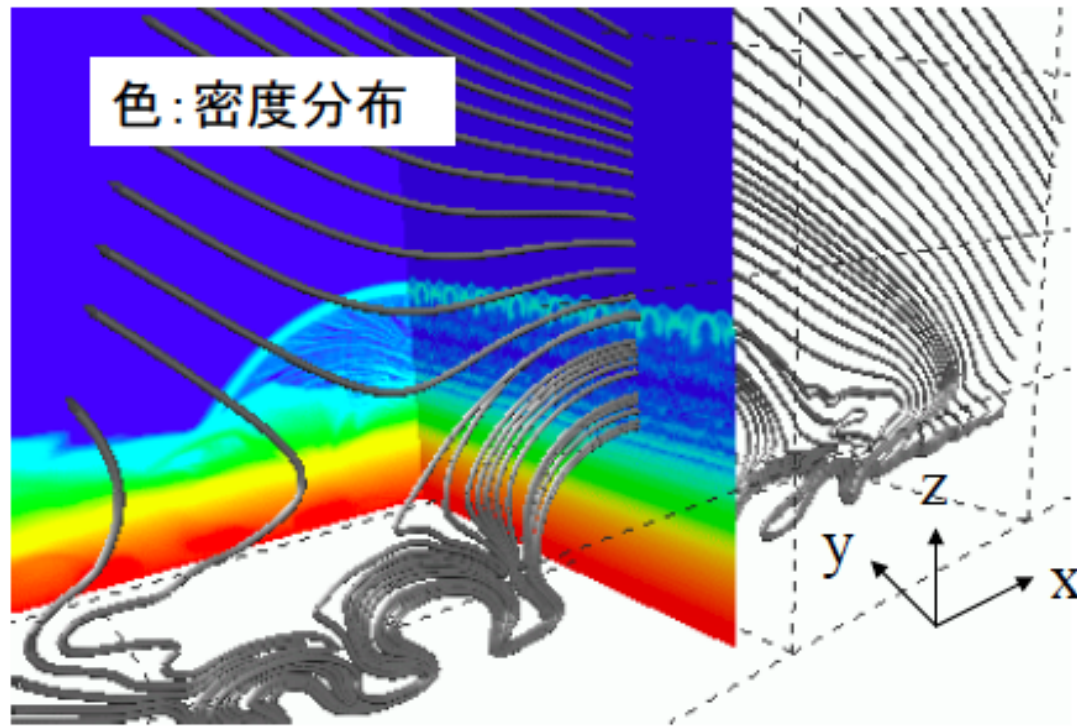
local current sheet thinning



Petschek-like reconnection (Isobe in prep.)

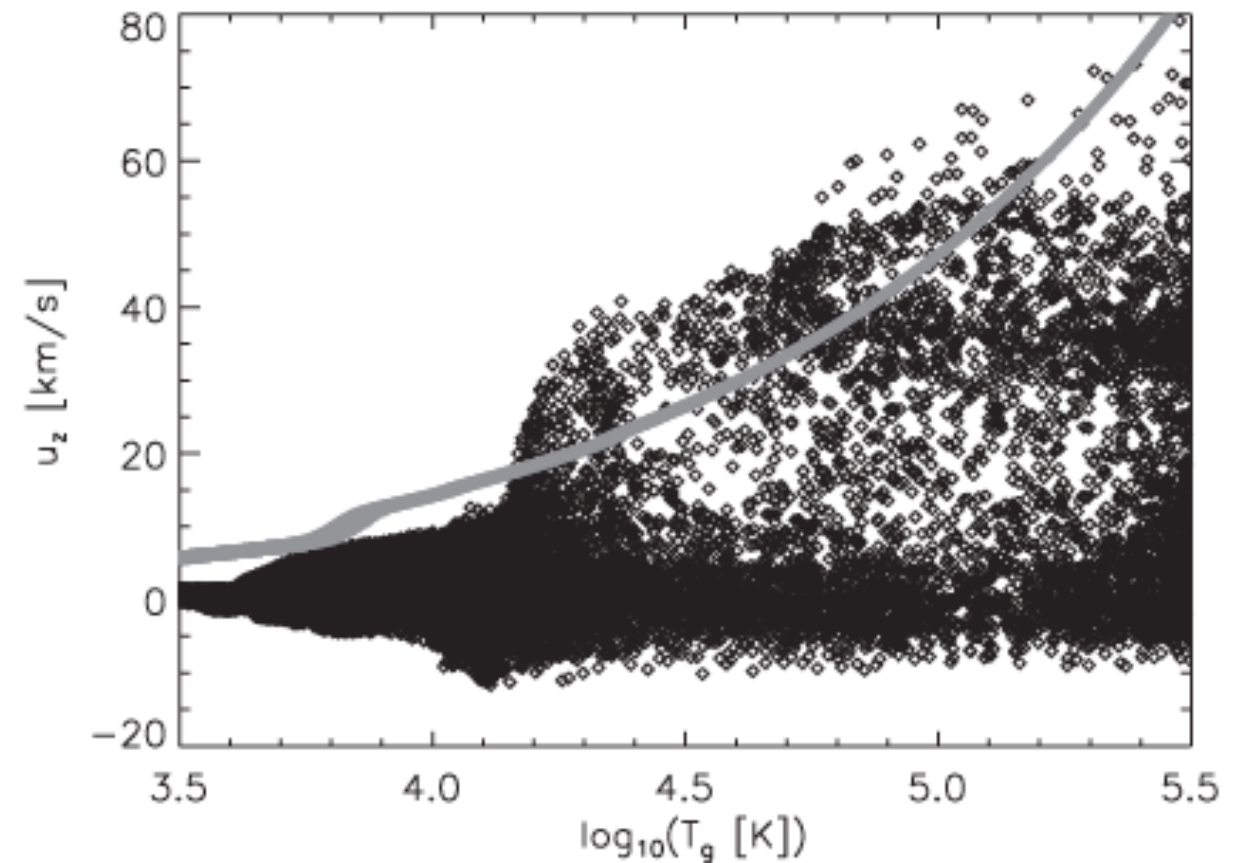
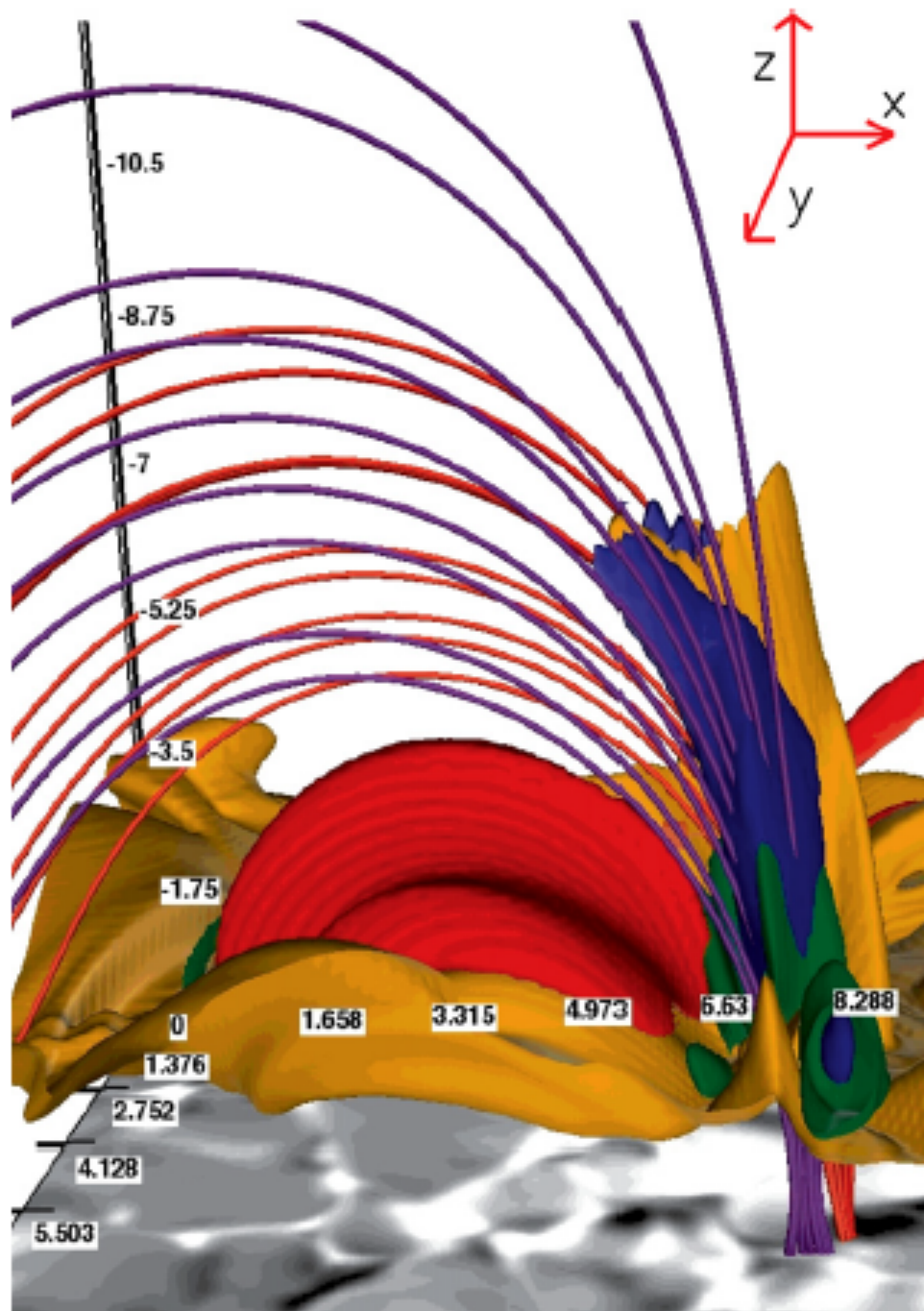


# Future Work: 3D Effects and Partially Ionization Effects



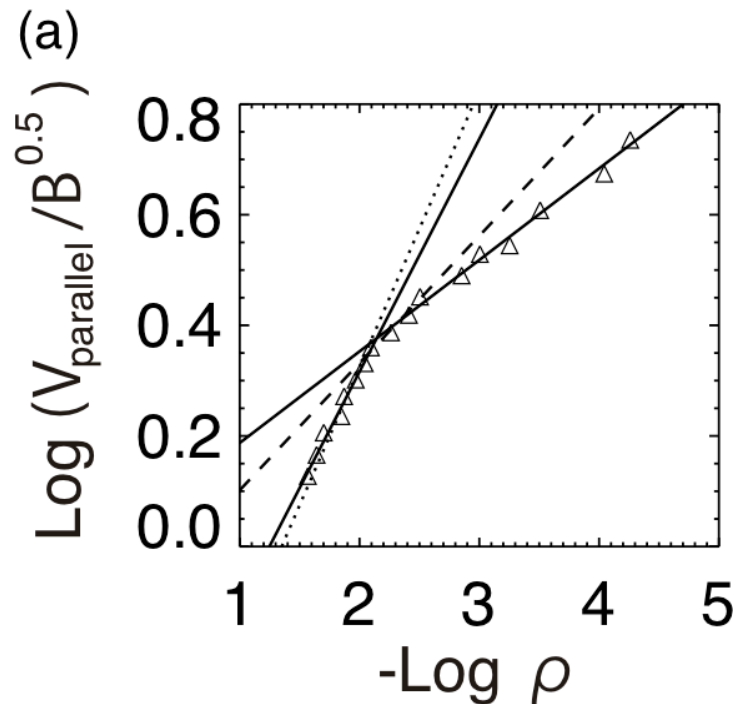


# Squeezing Acceleration



Results of the squeezing acceleration by the Lorentz force (Martinez-Sykora + 2012). This cannot account for jets which is much faster than the sound speed.

# Comparison of 2D results with 1D analytical relations

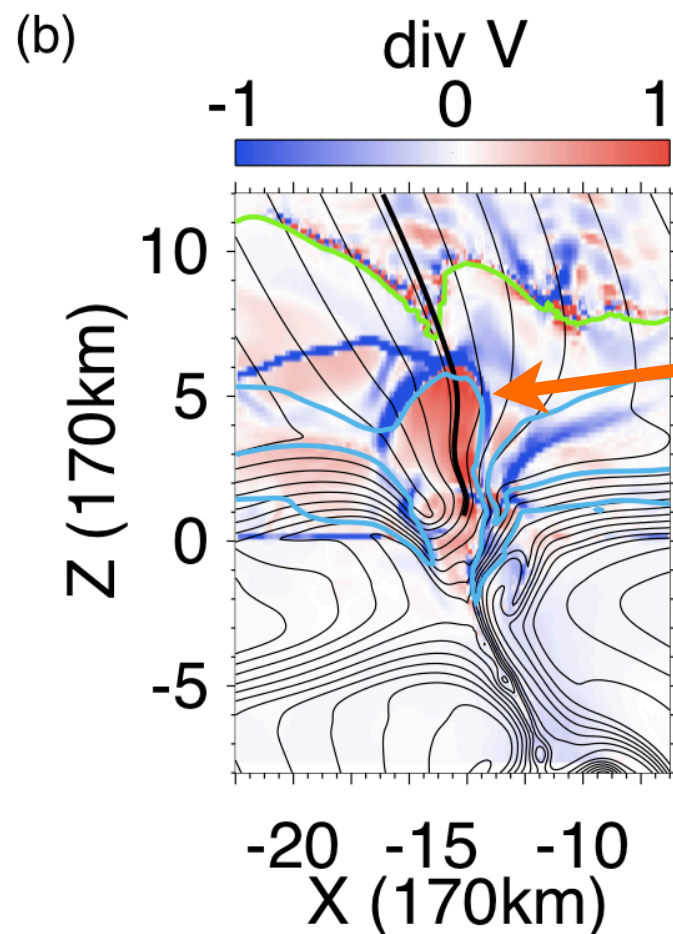


1D analytical relations

$$V_{parallel} B^{-0.5} \propto \begin{cases} \rho^{-0.5} & \text{(linear wave)} \\ \rho^{-0.236} & \text{(strong shock)} \end{cases}$$

2D simulation results

$$V_{parallel} B^{-0.5} \propto \begin{cases} \rho^{-0.36} \\ \rho^{-0.15} \end{cases}$$

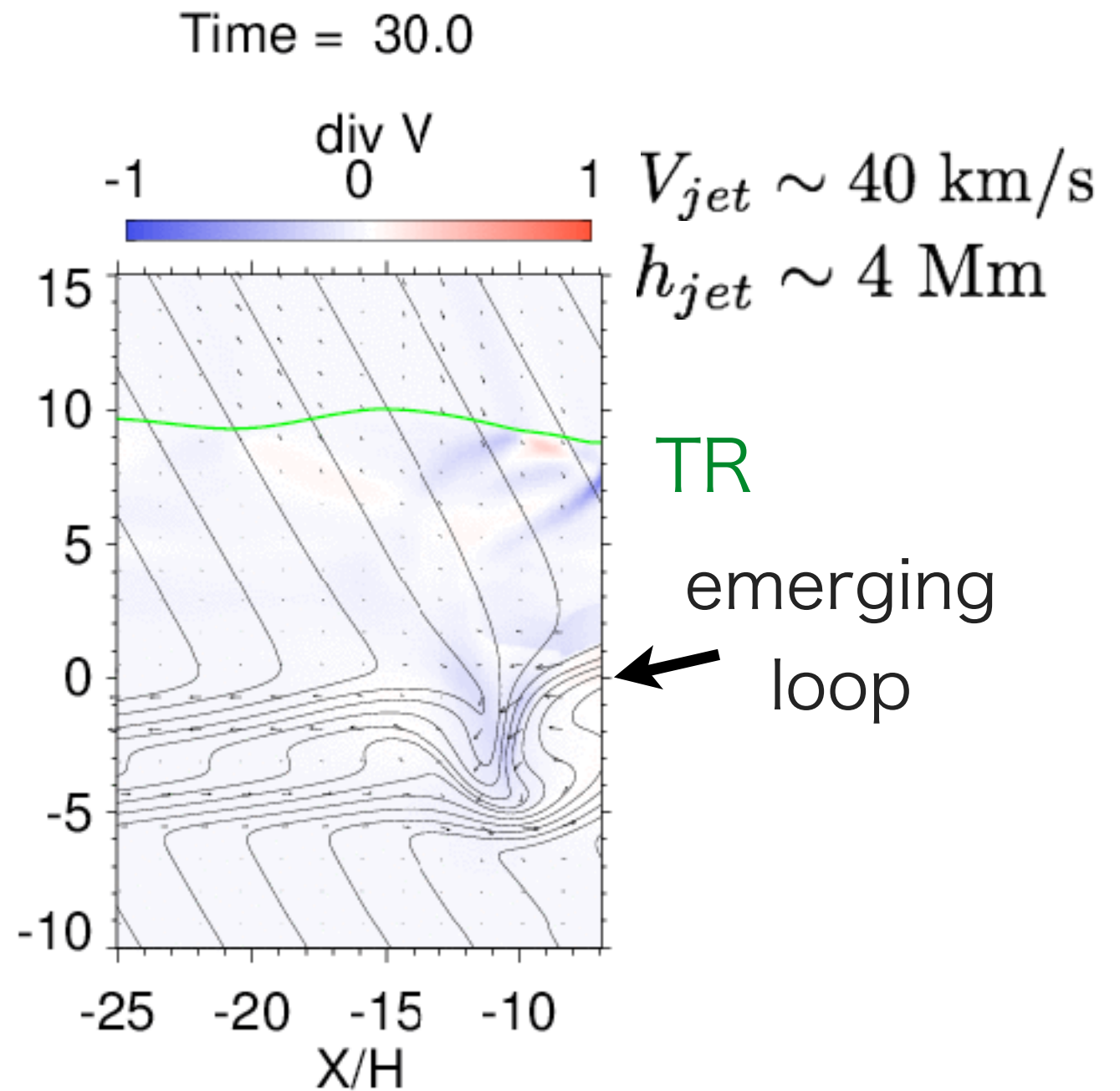
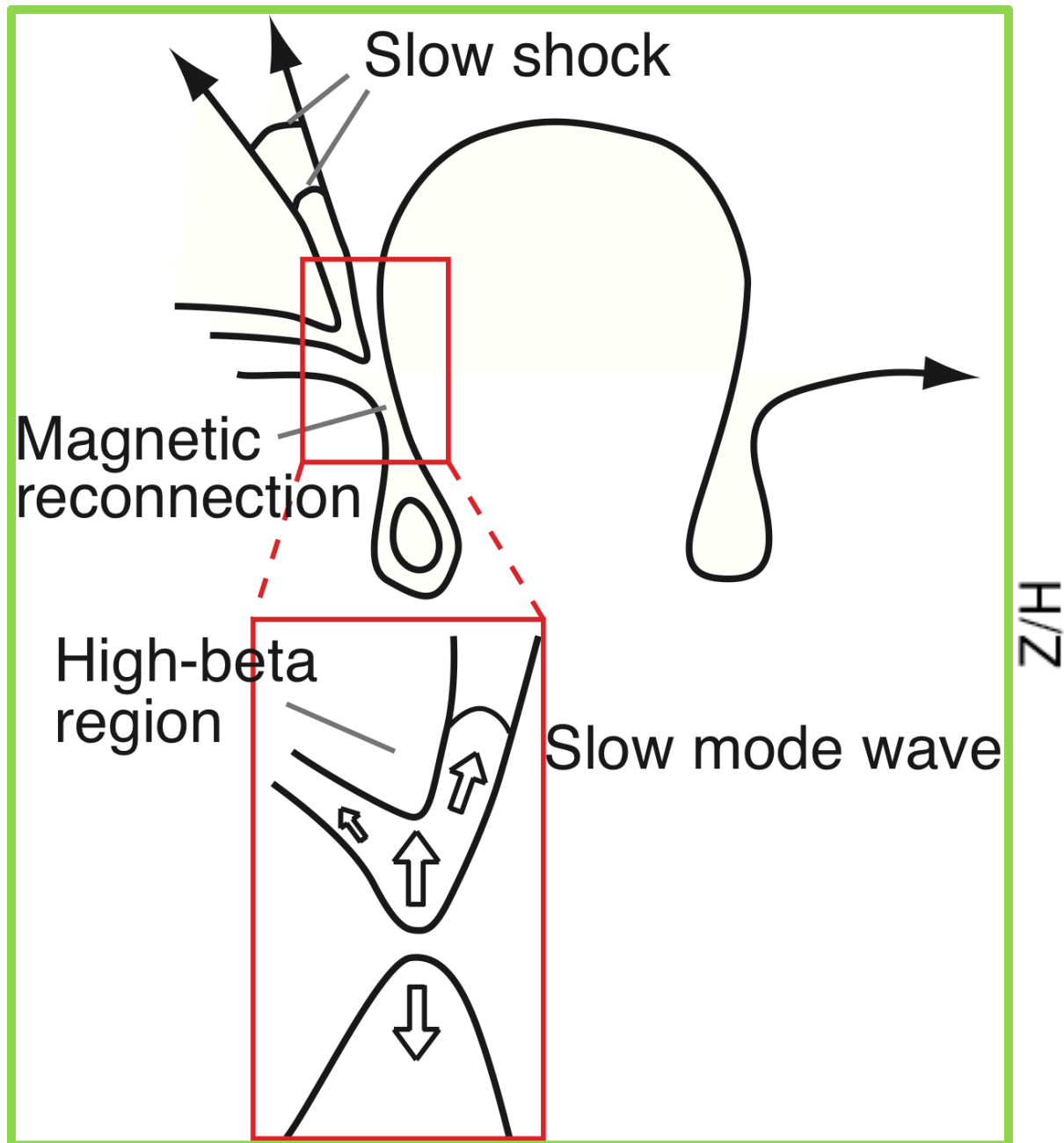


blue line: beta = 1

Plasma beta behind the shock is close to unity! Therefore the rigid flux tube approximation is broken even in the upper chromosphere.



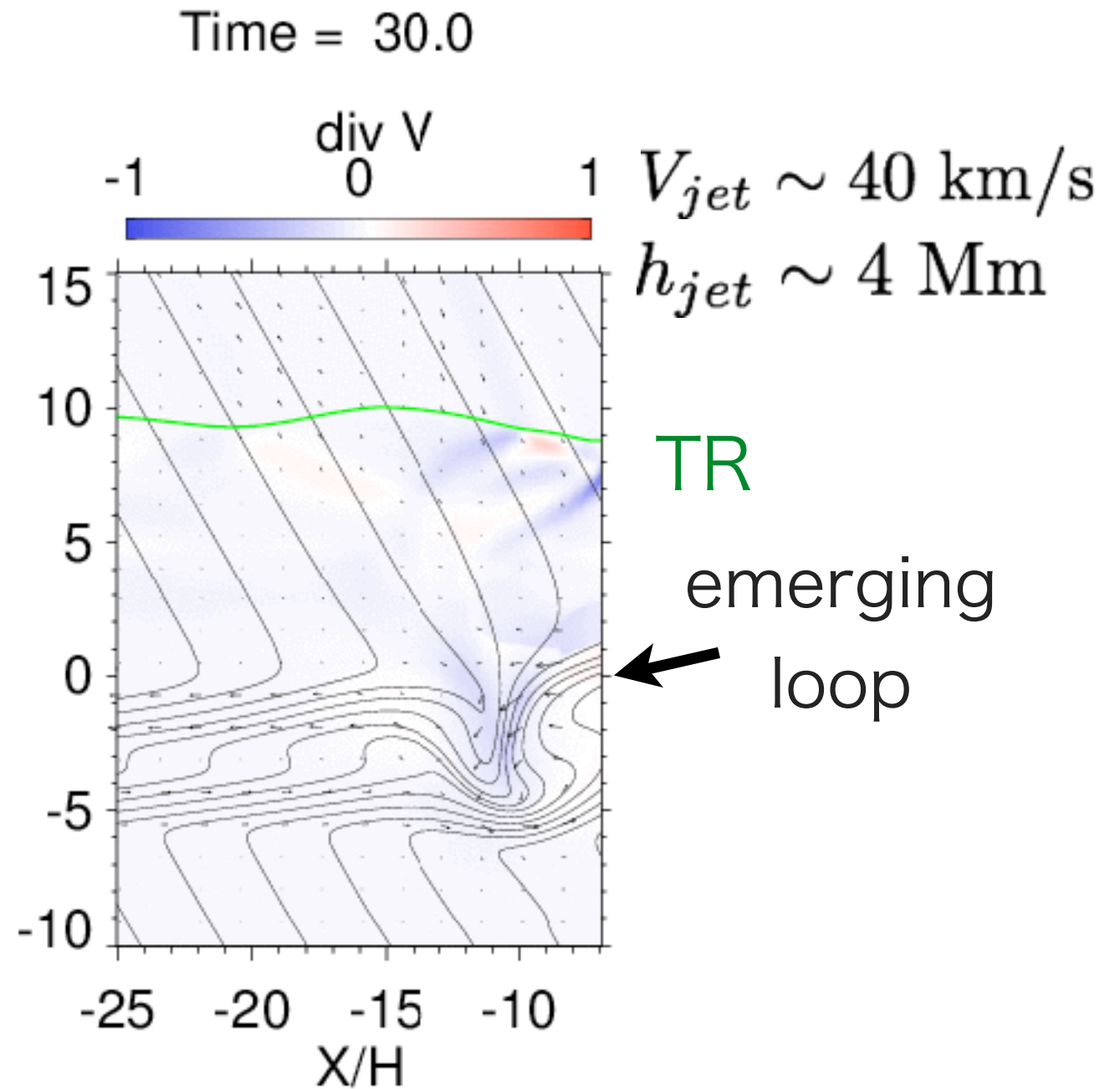
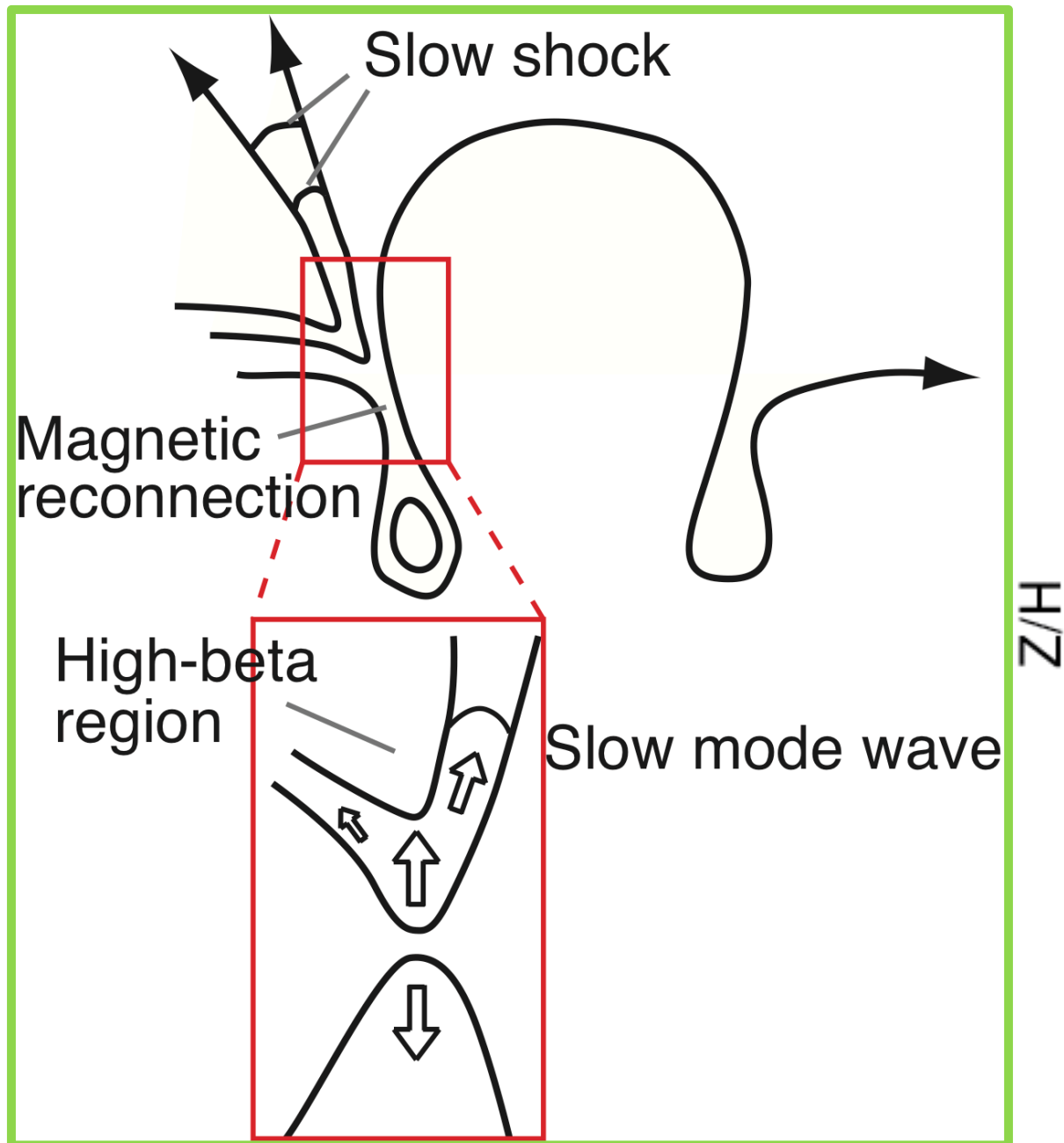
# Lower Atmospheric Recon.: Shock Acceleration



Note: The plasma behind the shock is strongly expanded.

This could be important for decreasing the optical depth,  
so disappearance of jets

# Lower Atmospheric Recon.: Shock Acceleration



Application:

Ellerman bombs  $\Rightarrow$  H-alpha Surges (e.g. Pariat+2004)

# Systematic Understanding of Chromospheric Jets: Classification by the Height of Recon. Points

Height of Reconnection Point

Lower/Middle Chrom.

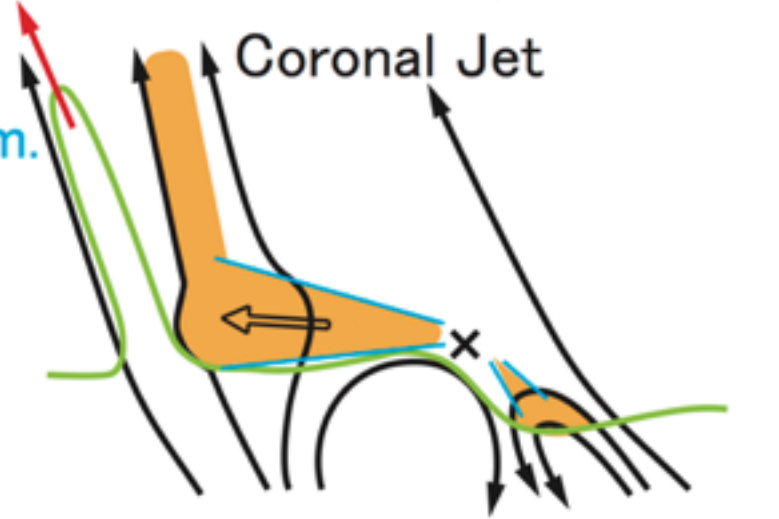
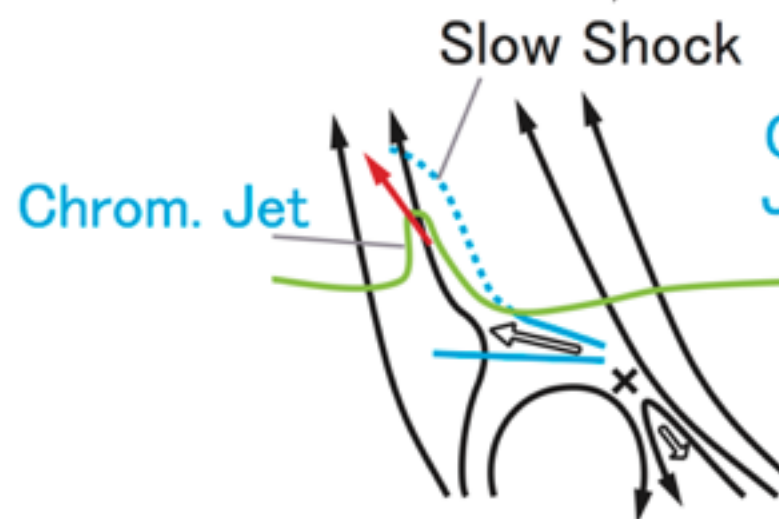
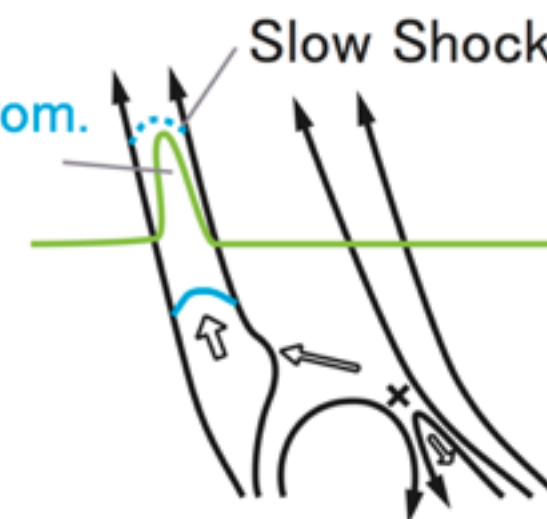
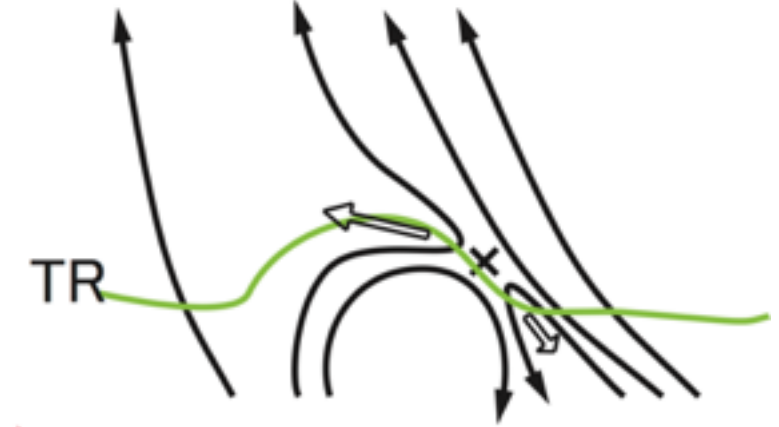
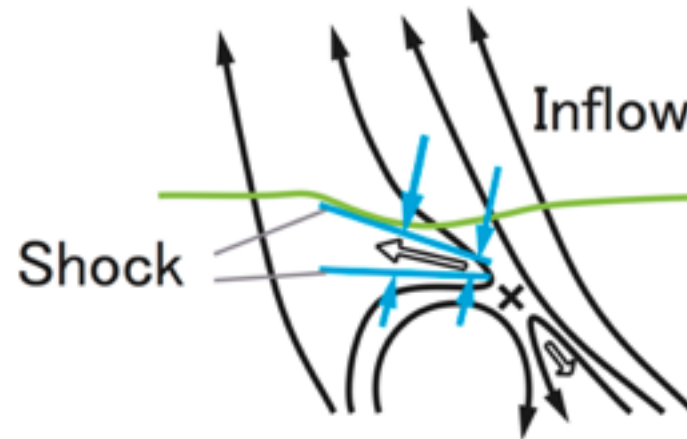
Upper Chrom.

TR/Corona →

Shock Acceleration

Shock + Whip-like Acceleration

Whip-like Acceleration  
(Yokoyama & Shibata 1996)



x : Reconnection Point  
⇨ : Reconnection Outflow

Ellerman bombs => surges  
tiny Ca jets  
some fraction of spicules?

larger Ca jets  
(~2Mm)

H-alpha surges