







1.3 Numerical study of Bz temperoal Behavior during 1997 Junuary Event(Sci in China)

In space weather event studies, the prediction of IMF Bz is very important. Using solar observation as inputs, the propagating process of January 1997 event is numerically simulated. The temporal Behavior of IMF Bz near the earth is shown to be in agreement with WIND data.







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1.5 MHD Shock Jump Condition (Solar Physics)

•MHD shock jump conditions have been used for more than 50 years after Ruderman proposed in 1950; We propose a new jump shock jump

condition by introducing the shock manifold and combining jump condition of the normal derivatives and the plain jump conditions. In this way, a multi-step method for explosive shock waves in solar wind can be established and extend some classic results of shock theory:

The method provides a way of predicting the arrival of shock near the earth. 2004.09 Beijing

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 $\frac{1}{(l-u_0)^2}\frac{d(l-u_0)}{dt} + \frac{\frac{l'-1}{l'+1} + \frac{u_0}{l'+u_0}}{1 + \sqrt{\frac{l'}{l'-u_0}}}\frac{1}{l'} = 0$

which is much simpler than (NeT). In this case, we compute smally by preconding as done in the above ments Benicken, of $\alpha_n \rightarrow 0$ in (27), then we can obtain a p for shock values I and in transitions R(n).

 $\delta \sim \exp\left[-\frac{\gamma+1}{(1+\gamma)(1+\sqrt{\frac{\gamma}{(1+\gamma)}})}\ln \theta\right]$

fach in considered by Wei (1982). On the other hand, if $e_k = 0$, considering the lecting the constitution of the gravity, i.e., the lo lecting the constitution of the gravity, i.e., the lo lecting the constitution of the gravity (1974) and

 $\left(\sqrt{\frac{2p}{p-1}}+2\right)\left|\frac{dt}{dt}+\frac{1}{p}\frac{dp}{dt}-1\right.$

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- 2, Numerical Schemes for the construction of MHD models
 - 1) No free parameter, Non-oscillatory, Dissipative scheme (NND—A TVD Type)

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- 2) Modified LAX-FRIEDRICHS-TVD scheme
- 3) Mixed GLM-MHD formulation method

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2.1 Initial Boundary inputs

How to use limited solar observation as inputs to initialize the solar wind simulation in order to obtain realistic steady state solar wind is an important step in solar wind structure simulation. This kind of study can provide a realistic solar wind background for numerical space weather model.

In order to give a realistic initial boundary condition for numerical steady state solar wind model, a solar wind background model constrained to solar observations from WSO Magnetic observation and HAO K-coronal brightness are established such that three dimensional magnetic field, density and velocity can be properly obtained.

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2.4 Mixed Generalized-Lagrange-Multiplier (GLM)-MHD formulation method

A new numerical scheme of 3rd order Weighted Essentially Non-Oscillatory (WENO) type for 2.5D mixed GLM-MHD in Cartesian coordinates is proposed. To show the validation and capacity of its application to MHD problem modelling, interaction between a magnetosonic shock and a denser cloud and magnetic reconnection problems are used to verify this new MHD code. It is significant to note that a new divergence cleaning approach for the MHD equations by coupling the divergence constraint with the evolution equations using a generalized Lagrange multiplier (GLM) is employed. The mixed hyperbolic/parabolic GLM ansatz offers both propagation and damping of divergence errors. Moreover, the magnetohydrodynamic part of the GLM-MHD system is still in conservation form by taking account of the arguments by Dellar and Dedner.

$$\frac{\partial}{\partial t} \begin{pmatrix} \rho \\ \rho \mathbf{v} \\ \mathbf{B} \\ \psi \\ \psi \end{pmatrix} + \left\{ \nabla \cdot \begin{pmatrix} \rho \mathbf{v} \mathbf{v} + (p + \frac{1}{2} \mathbf{B} \cdot \mathbf{B}) \mathbf{I} - \mathbf{B} \mathbf{B} \\ \mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v} + \psi \mathbf{I} \\ \mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v} + (\mathbf{v} \cdot \mathbf{B}) \mathbf{B} \\ \mathbf{c}_{\mathbf{Z}}^{2} \mathbf{B} \end{pmatrix} \right)^{T} \right\}^{T} = -\nabla \cdot \mathbf{B} \begin{bmatrix} 0 \\ 0 \\ \mathbf{v} \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ -\mathbf{B} \cdot (\nabla \psi) \\ -\mathbf{u} \cdot (\nabla \psi) - \frac{2}{\epsilon_{\mathbf{J}}^{2}} \psi \end{bmatrix}$$

$$(U + p + \frac{1}{2} \mathbf{B} - \mathbf{B} \mathbf{v} - (\mathbf{v} \cdot \mathbf{B}) \mathbf{B} \\ \mathbf{c}_{\mathbf{Z}}^{2} \mathbf{B} \end{pmatrix}^{T} = -\nabla \cdot \mathbf{B} \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ -\mathbf{B} \cdot (\nabla \psi) \\ -\mathbf{u} \cdot (\nabla \psi) - \frac{2}{\epsilon_{\mathbf{J}}^{2}} \psi \end{bmatrix}$$

$$(Z + \mathbf{D} \cdot \mathbf{A} - \mathbf{D} \cdot \mathbf{A} - \mathbf{A} + \mathbf{A} - \mathbf{A}$$









3、Stu	idies of fundamental physical mechanisms between Sun-Earth system	n
3.1	Boundary Layer of IMC (J.G.R. G.R. L)	
	Based on a statistical analysis of the boundary physical states of 80 magnetic clouds reported in the literature from the years 1969 to 2001, we suggest a new identification the magnetic cloud boundary by describing it as front and tail boundary layers (BLs) formed through the interaction between the magnetic cloud and the ambient medium.	of
	The outer boundary of the layer often displays the properties of magnetic reconnection which could be characterized by a "three-high state" (relatively high proton temperature, high proton density, and high plasma beta) and the corresponding magne signatures (the magnetic intensity drop and the directions change abruptly; the abrupt azimuthal changes, delta(phi)–180, and latitudinal changes, delta(theta)–90, in the magnetic field).	, tic
	The inner boundary of the layer exhibits a "three-low state" (relatively low stratom temperature, two proton tensity, and low plasma beta) and separates the magnetic cloud tests, which has not basically been affected by the interactions, from the boundary layers.	
	The treat boundary layer could be associated with the outer loops of CMEs and its average time scale is 1.7 hours; the tail boundary layer seems not be a filament and its average time scale is 3.1 hours. The distribution function of magnetic fluctuations in boundary layer is significantly different from those in the ambient solar wind and the cloud body itself. The preliminary numerical simulation in principle confirms this new identification and could qualitatively explain most of the observations of the cloud boundary.	ne 7
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Examples: Boundary layer of interplanetary magnetic cloud (M_p, G_p) and shock, sheath and magnetic cloud body (let)

Magnetic cloud boundary layer (right)

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3.2 possible mechanism of MC boundary layers(I.G.R.) MC boundary layers may be produced by magnetic reconnection. This mechanism can explain the basic observations of MC and is principally produced by numerical simulation. At the outer boundaries, magnetic strength decreases abruptly, its directions also changes abruptly(△0-180 * .△0 ≥45*) and have three high sibility, high temperatures, high plasma beta, at the inner boundary, magnetic strength increases, its directions change gradually and have three low states;; low density, low temperature and low plasma beta.





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3.4 Global structure of coronal mass flux output and magnetic field (J. G. R) The basic characteristics of the coronal mass output near the Sun are analyzed with the statistic and numerical methods by using observational data from K corona brightness, IPS(interplanetary scintillation), and photospheric magnetic field during the descending phases (1983) and the minimum (1984) of solar activity from the point of view of the global distribution of the solar magnetic field on the source surface (at 2.5 solar radii (Rs)).

The main results are as follows: (1) There are certain regular persistent patterns in the global distributions of coronal mass outputs flux Fm (density x speed V), which shows that the highest Fm in 1983 and 1984 display more regularly double peaks and single-peak wavelike patterns on the source surface (2.5 RS), respectively. The highest and the lowest Fm are associated with the coronal current sheet and the polar corona regions, respectively, and the other regions are associated with a moderate Fm. (2) The speed dependence of Fm is different for various magnetic structures.

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By solving a self-consistent MHD system based on the observations of K coronal brightness and the photospheric magnetic fields, a

preliminary numerical study of the global distribution near 2.5 Rs for the Carrington rotation 1742 in 1983 has been made.



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Residence of the Carrington rotation 1942 in 1983. They are obtained from the MHD model with the observational data of the photospheric magnetic fields and K coronal brightness as inputs

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- 4. Solar disturbances & geomagnetic storm prediction:
- statistical methods 4.1 Geomagnetic disturbance ISF prediction method(Adv Space Res, Sci in China)



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4.2 Solar wind parameters prediction-SSP method (JGR) Based on K-coronal brightness and photospheric magnetic field as inputs, a simplified MHD method is proposed to obtain the solar surface parameters at 2.5 Rs and then empirical relation is used to derive the solar wind parameters near the earth(Source Structure Prediction Method, SSP method).





